

# Modelling of electronic and transport properties in semiconductor nanowires

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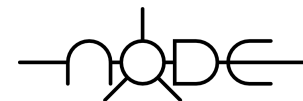
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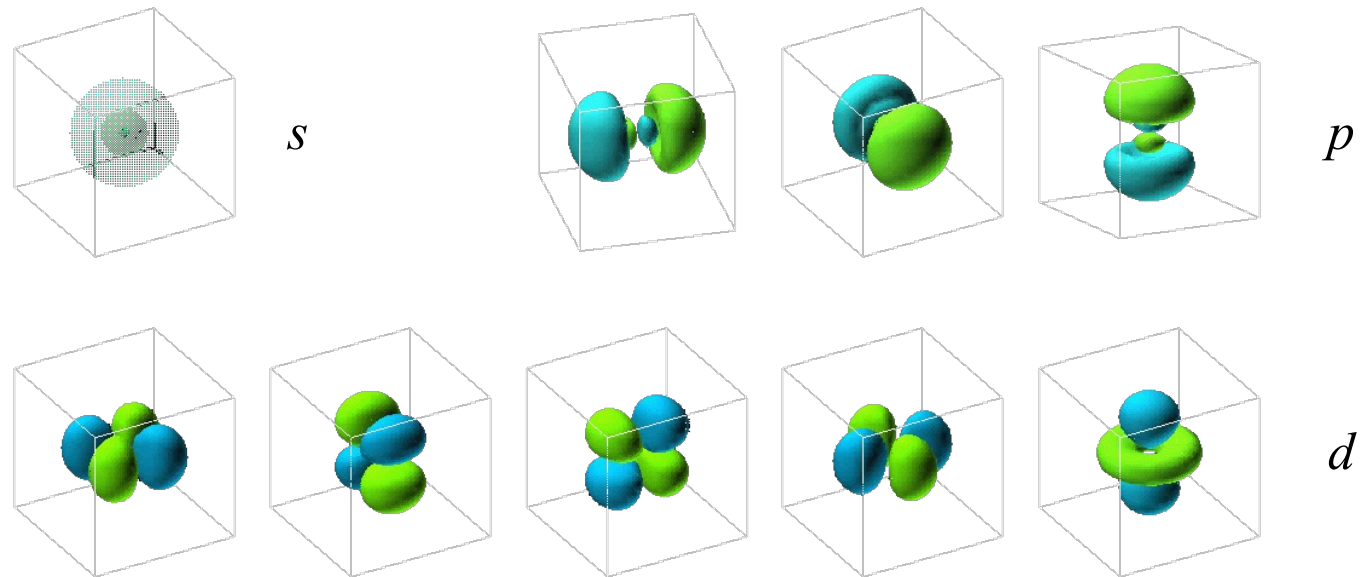


- NODE – Nanowire-based One Dimensional Electronics
- Bottom up approach to nanowire based electronics
- Evaluate and deliver replacement and add-on technologies to silicon CMOS
- Partners within **academia** and **industry** in Europe
  - Lund University (**Coordinator**)
  - TU Delft, Max-Planck Institute (Halle), Scuola Normale Superiore (Pisa)
  - Würzburg University, CEA, IBM Research, IMEC, Philips, Infineon, Qumat
- CEA Grenoble is contributing with **theory** and **structural analysis** of nanowires

- **Strains in nanowire heterostructures.**
  - Effects of strain relaxation of nanowire heterostructures.
  - Barrier lowering due to strain.
- Doping of nanowires
  - Dielectric effects
  - Increased binding energies
- **Transport properties of semiconductor nanowires.**
  - The Kubo-Greenwood and Landauer-Büttiker approach.
  - Application : Surface disorder.



# The tight-binding method



- Principle : Expand the wavefunctions as linear combination of atomic orbitals.
  - The **range** of the model is limited to first, second or third-nearest neighbors.
  - The matrix elements of the hamiltonian are considered as adjustable parameters usually fitted on bulk band structures then **transferred** to the nanostructures.
  - **The computation time scales (at least linearly) with the number of atoms** (up to a few millions of atoms today).

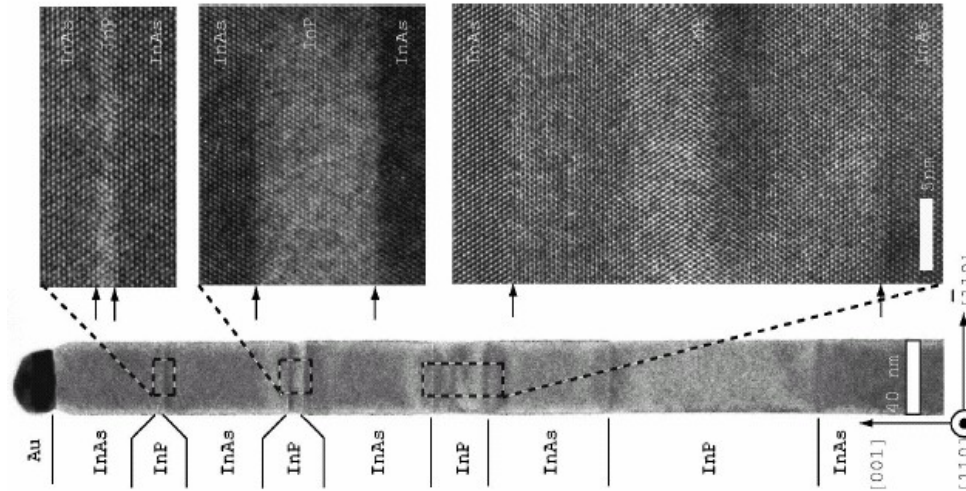
## Part I :

# Strains in nanowire heterostructures

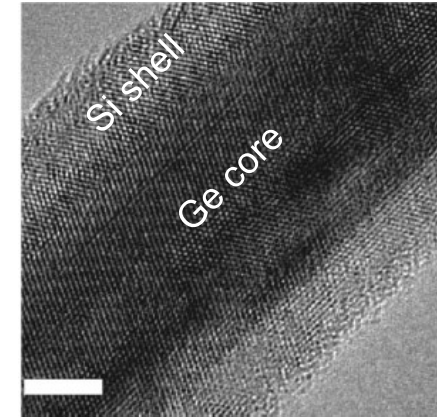
# Nanowire heterostructures

- Large interest in nanowire « heterostructures » for optics & transport :

M. T. Björk *et al.*, Appl. Phys. Lett. **80**, 1058 (2002).

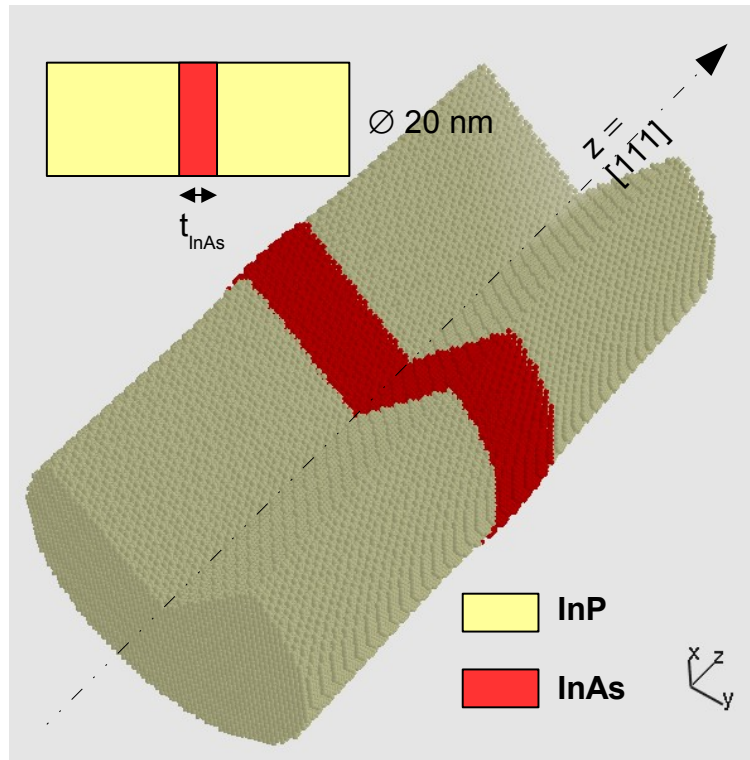


W. Lu *et al.*, PNAS. **102**, 10046 (2005).



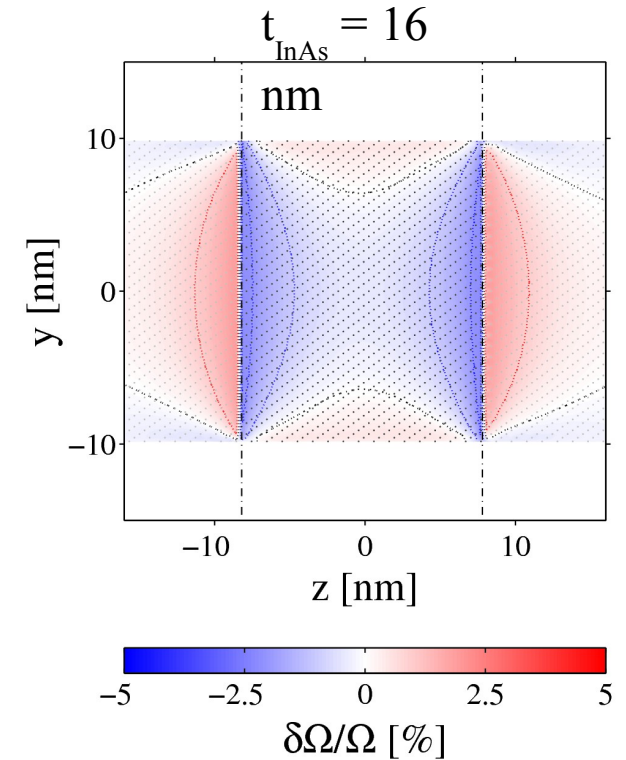
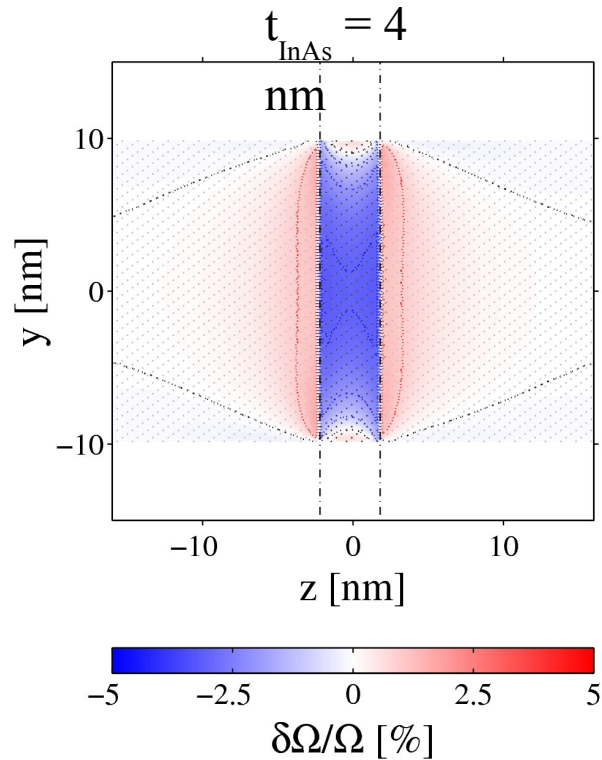
- Strain relaxation is believed to be efficient in these structures, likely allowing the growth of thick lattice mismatched layers.
- A few issues :
  - **What is the effect of strain relaxation on the electronic properties of nanowire heterostructures ?**
  - **What is the effect of an overgrown shell ?**

# InAs/InP heterostructures



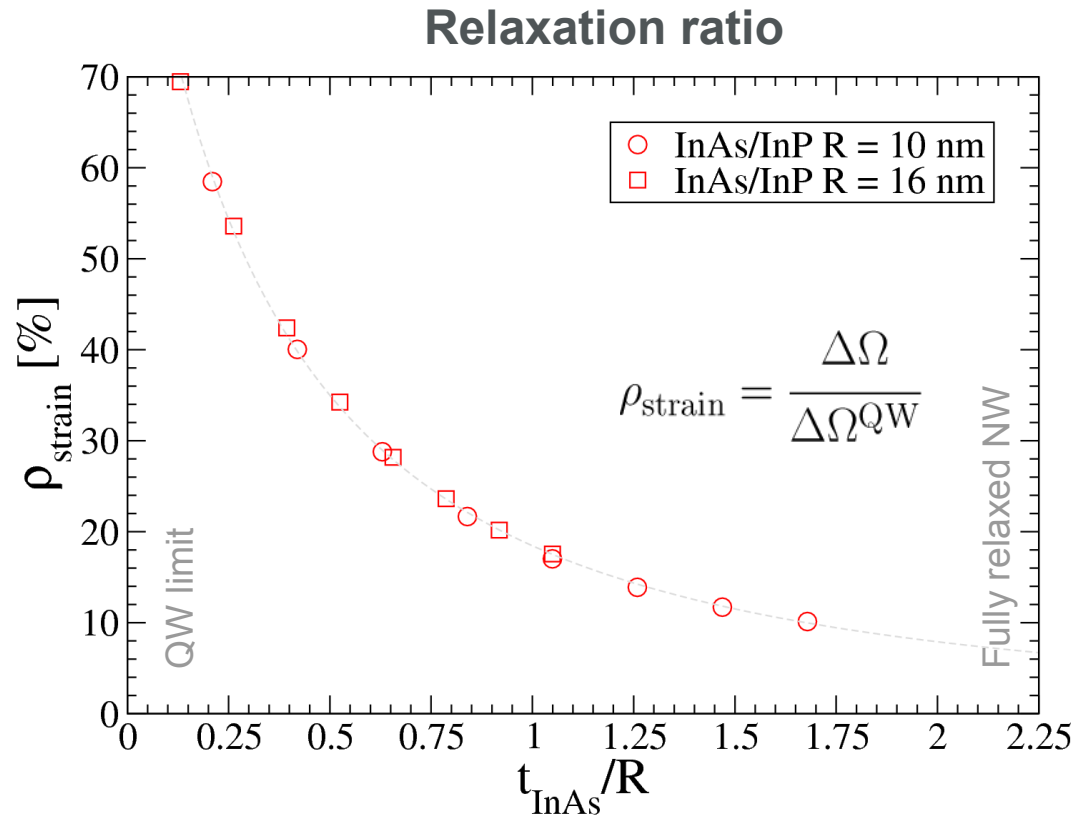
- The bond length is 3.13% shorter in InP than in InAs. The InAs layer is thus **compressed** by the InP core, but can partly **relax** strains at the surface of the nanowire.
- Strain relaxation is computed using Keating's Valence Force Field model

# Strain relaxation



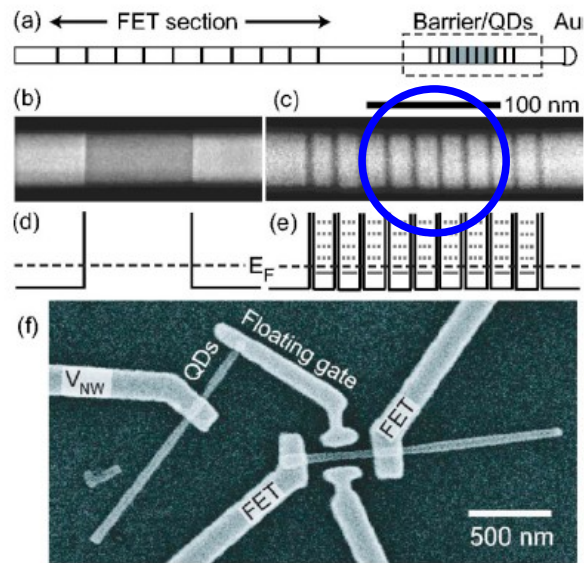
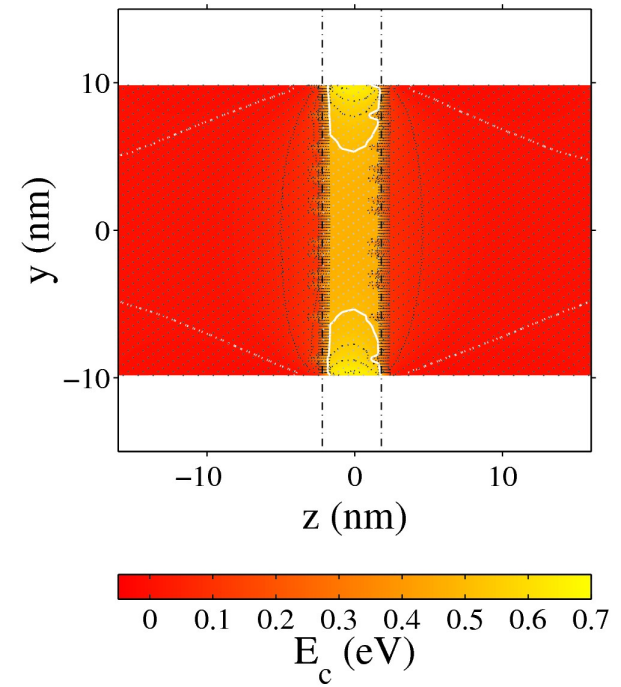
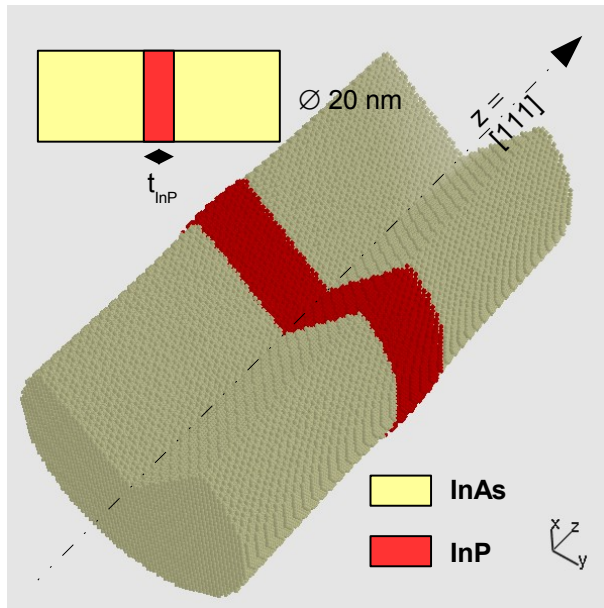
- Strain relaxation is very **efficient** in nanowire heterostructures. The InAs layer expands outwards and distorts the surface of the nanowire. The strain distribution is however very **inhomogeneous** in thin InAs layers : the surface is overrelaxed while the axis is still significantly compressed.





- The InAs layers are almost **completely relaxed** when  $t_{\text{InAs}} > 2R$ .

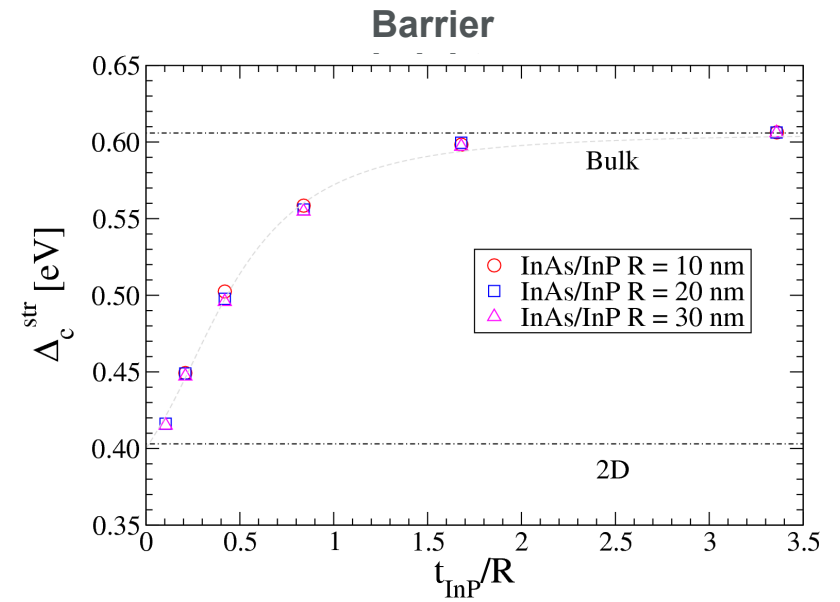
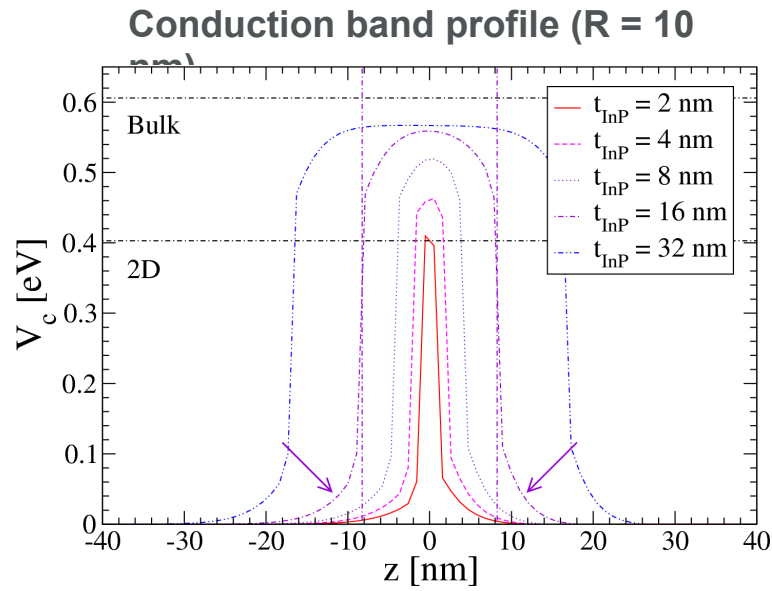
# InP tunnel barriers in InAs nanowires



H. A. Nilsson *et al.*, *Applied Phys. Lett.* **89**, 163101 (2006).

- The InP barrier is dilated by the InAs core, which tends to **lower the conduction band energy**.

# InP tunnel barriers in InAs nanowires



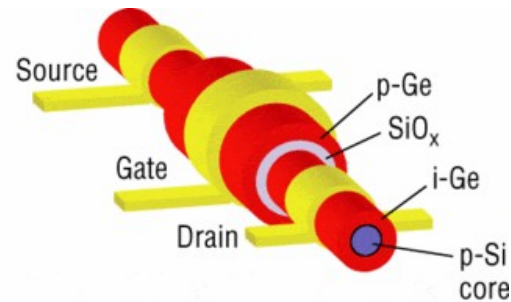
- The barrier height is **close to the bulk value (0.6 eV)** in thick InP layers ( $t_{\text{InP}} > 1.5R$ ), but tends to the **2D limit (0.4 eV)** in thin ones.

## Part II :

# Doping in nanowires

# Doping in nanowires

- Doping of semiconductor nanowires is an important issue and challenge.



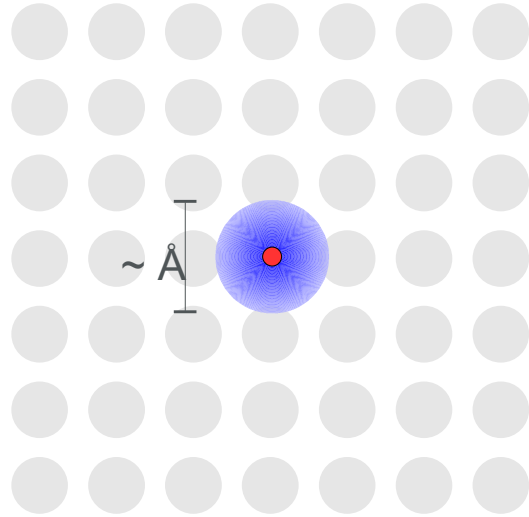
L. J. Lauhon *et al.*, Nature **420**, 57 (2002)

- Non uniform dielectric environments  $\Rightarrow$  Complex electrostatics.

**What is the binding energy of donor and acceptor impurities ?**  
**What is the doping efficiency ?**  
**How can we improve doping efficiency?**

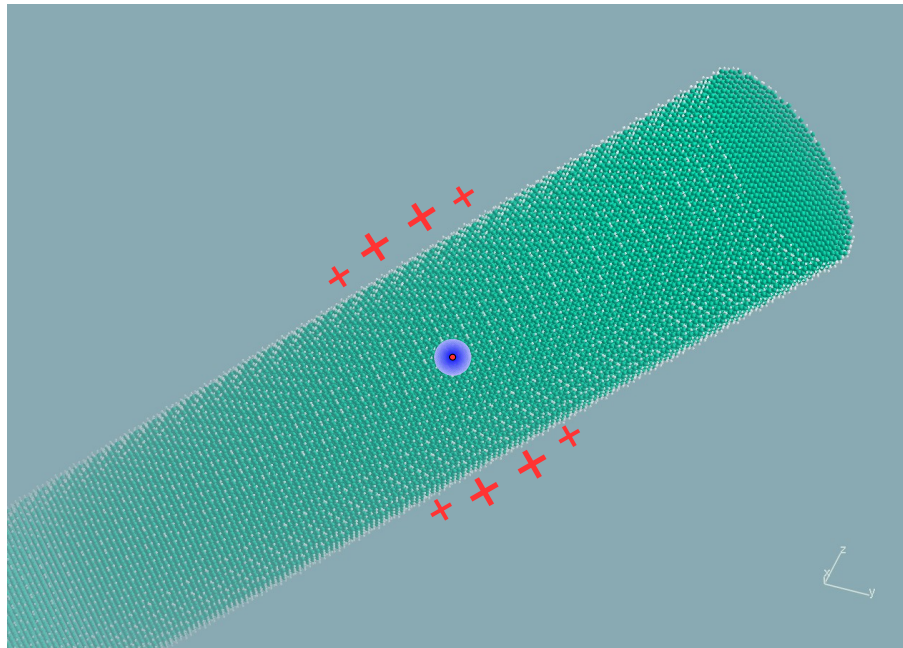
- Conclusions :
  - The binding energy of donor and acceptor impurities can be greatly enhanced in thin ( $R < 10$  nm) silicon nanowires in vacuum.
  - High- $\kappa$  dielectrics and all-around metallic gates for example can help to improve the doping efficiency.

# The microscopic interpretation of classical electrostatics



A diagrammatic equation enclosed in a dashed box. It shows a red dot followed by a plus sign, then a blue cloud, followed by an equals sign and another blue cloud. Below this, the corresponding mathematical terms are shown:  $+1$ ,  $-\left(1 - \frac{1}{\kappa}\right)$ , and  $\frac{1}{\kappa}$ .

- An ionized donor attracts nearby valence electrons and **gets screened by a short-range « cloud » of negative charges**.
- **The impurity and its cloud behave as a total charge  $1/\kappa$**  creating a potential  $V(\mathbf{r}, \mathbf{r}') = 1/\kappa |\mathbf{r} - \mathbf{r}'|$  at long distances.
- In bulk materials, the charge  $-\left(1 - \frac{1}{\kappa}\right)$  in this cloud comes « from infinity ».



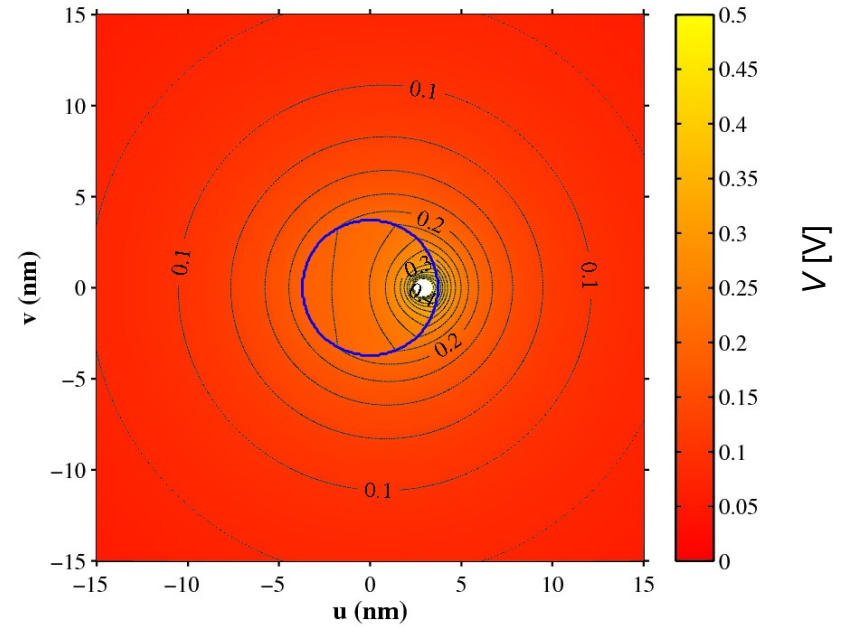
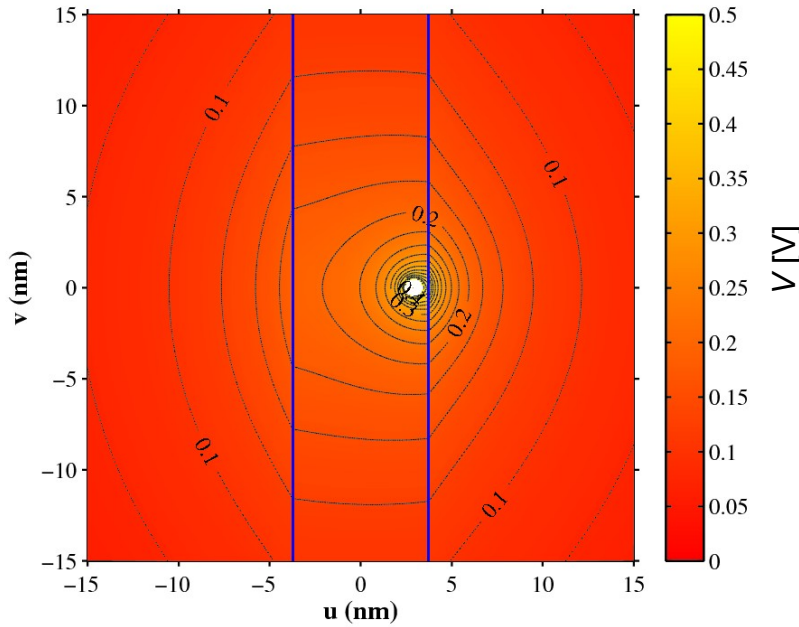
- In a nanowire, however, **the charge in the cloud comes from the surface** (« **image charges** » distribution).

- The solution of Poisson equation :

$$\nabla_{\mathbf{r}'} [\kappa(\mathbf{r}') \nabla_{\mathbf{r}'} V(\mathbf{r}, \mathbf{r}')] = 4\pi\delta(\mathbf{r} - \mathbf{r}')$$

is actually the potential created in vacuum by the (unscreened) impurity, its cloud and its image charges.

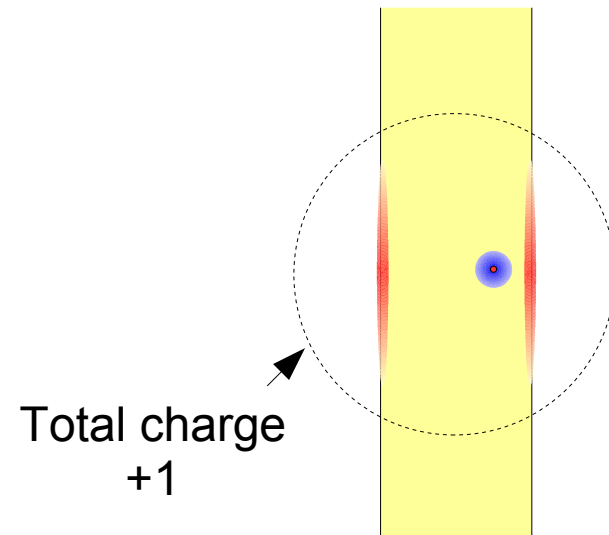
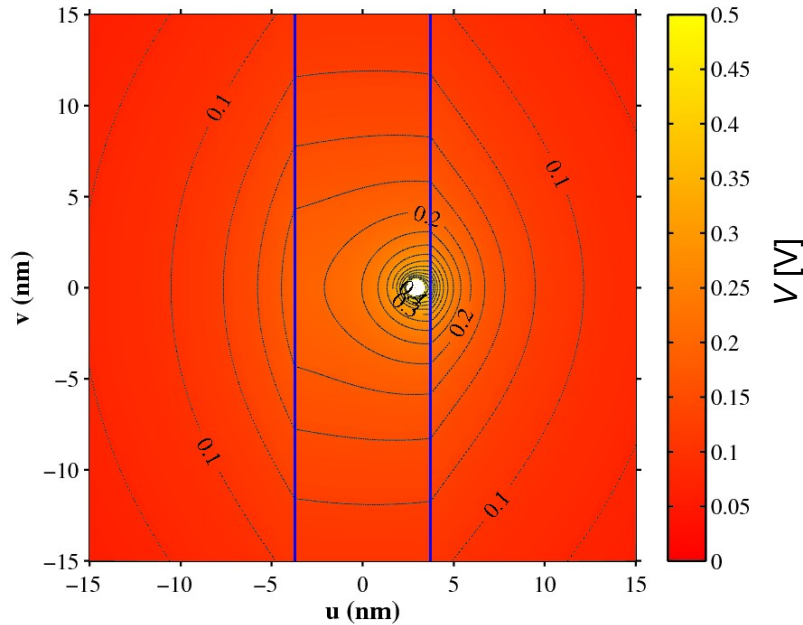
# The hydrogenoid impurity problem in nanowires



- The potential is not isotropic due to the image charges.

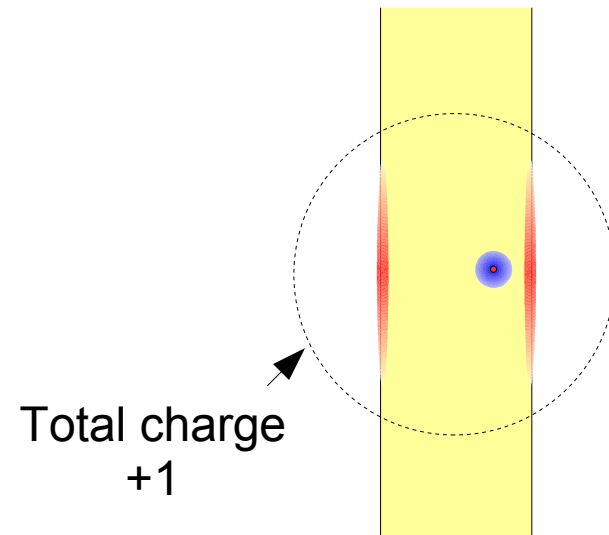
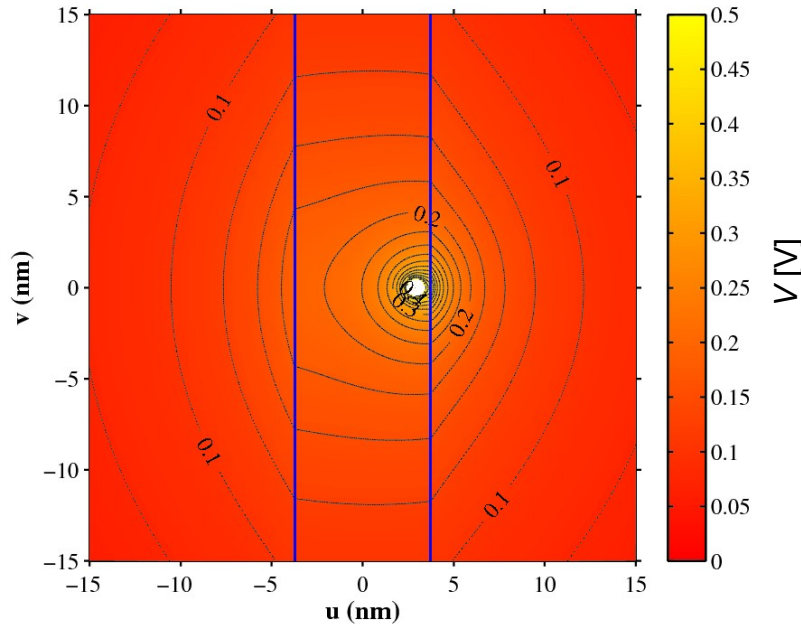


# The microscopic interpretation of classical electrostatics



- The potential is not isotropic due to the image charges.
- The total charge of the system (impurity + cloud + image charges) is +1 ; hence the potential decreases as  $1/|\mathbf{r} - \mathbf{r}'|$  far enough (a few  $R$ 's) from the impurity. As a consequence, **the potential around the impurity is deeper than in bulk.**

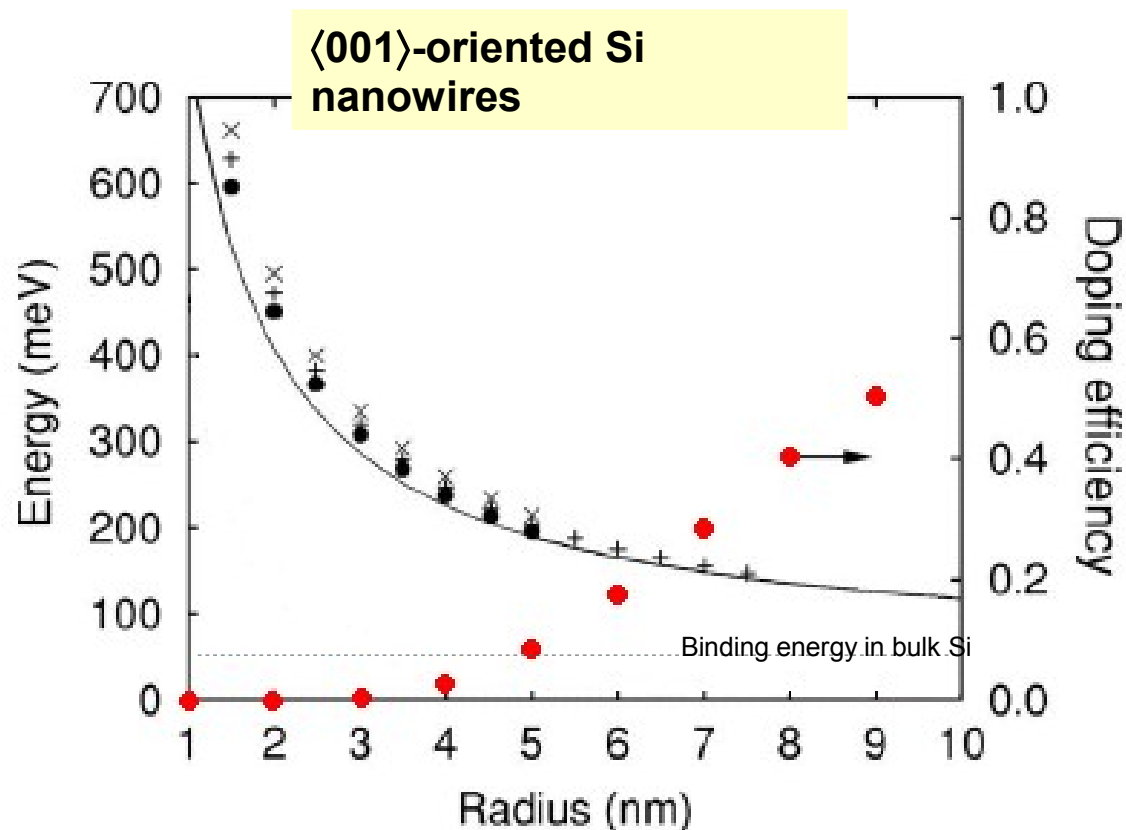
# The microscopic interpretation of classical electrostatics



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**Effects on the binding energies and doping efficiency ?**

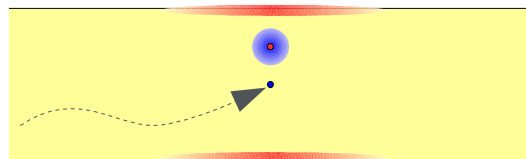
# Doping the nanowires



Binding energy of a donor in a Si nanowire as a function of its radius.

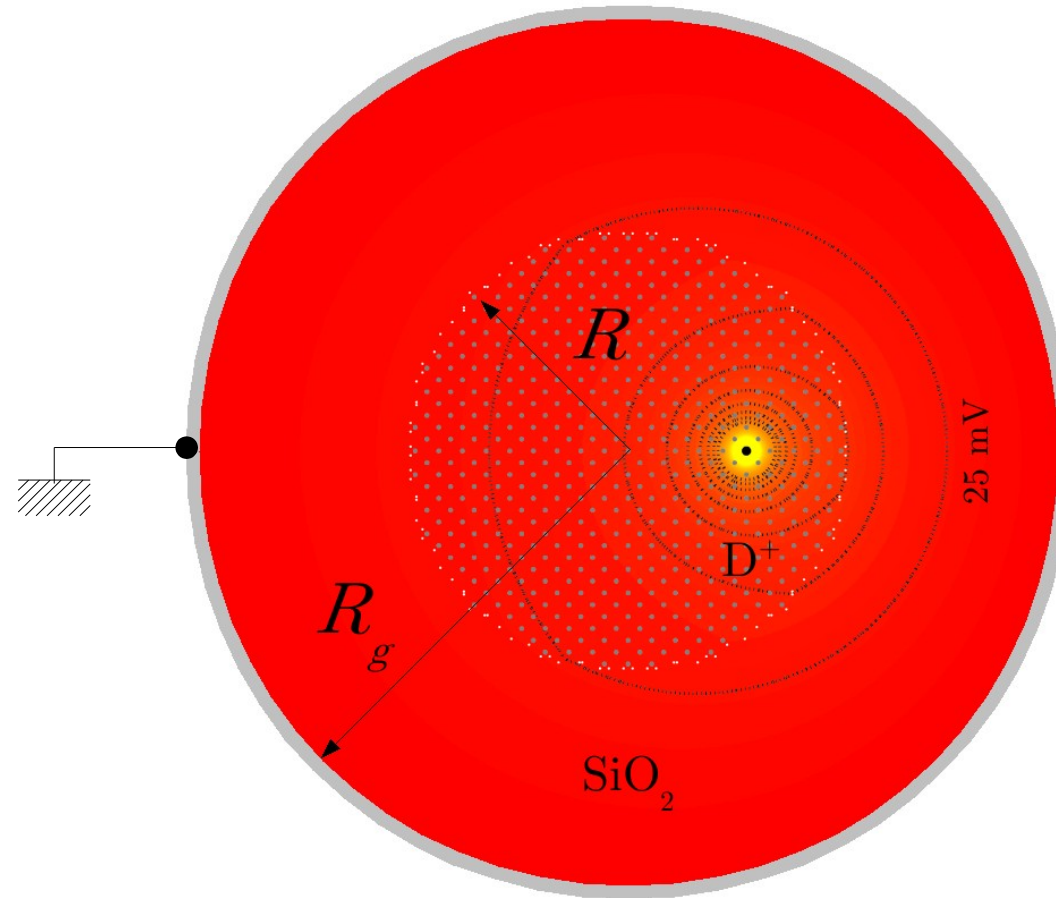
The donor is located along the nanowire axis.

- The image charges **increase the binding energy of the donor up to a few hundreds of meV** in the smallest nanowires !!



The electron is trapped around the donor by the impurity and its image charges.

# Screening in a complex dielectric environment

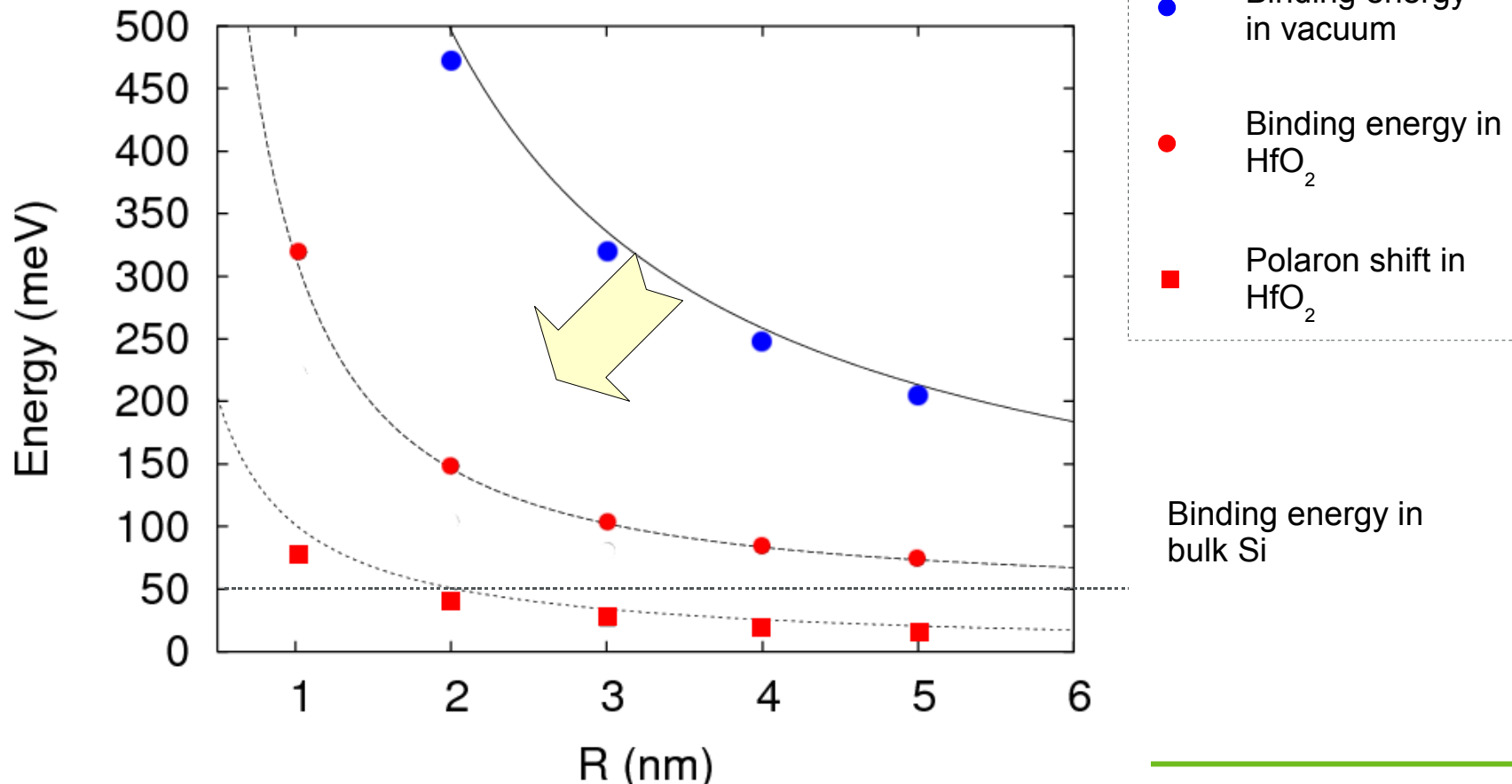


- Oxides and metallic gates screen the impurity potential...  
⇒ **Decrease of the binding energy**  
**... BUT ...**
- The dielectric response of the oxides is slow...  
⇒ **Polaronic enhancement of the binding energy !**

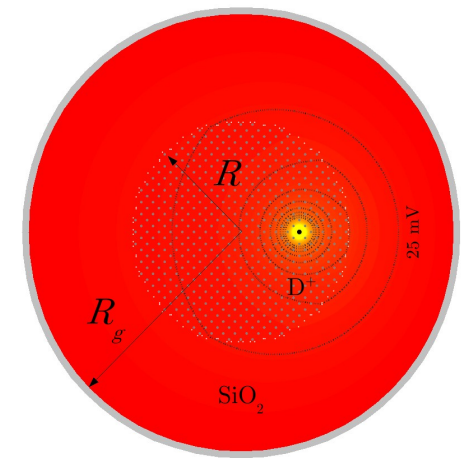
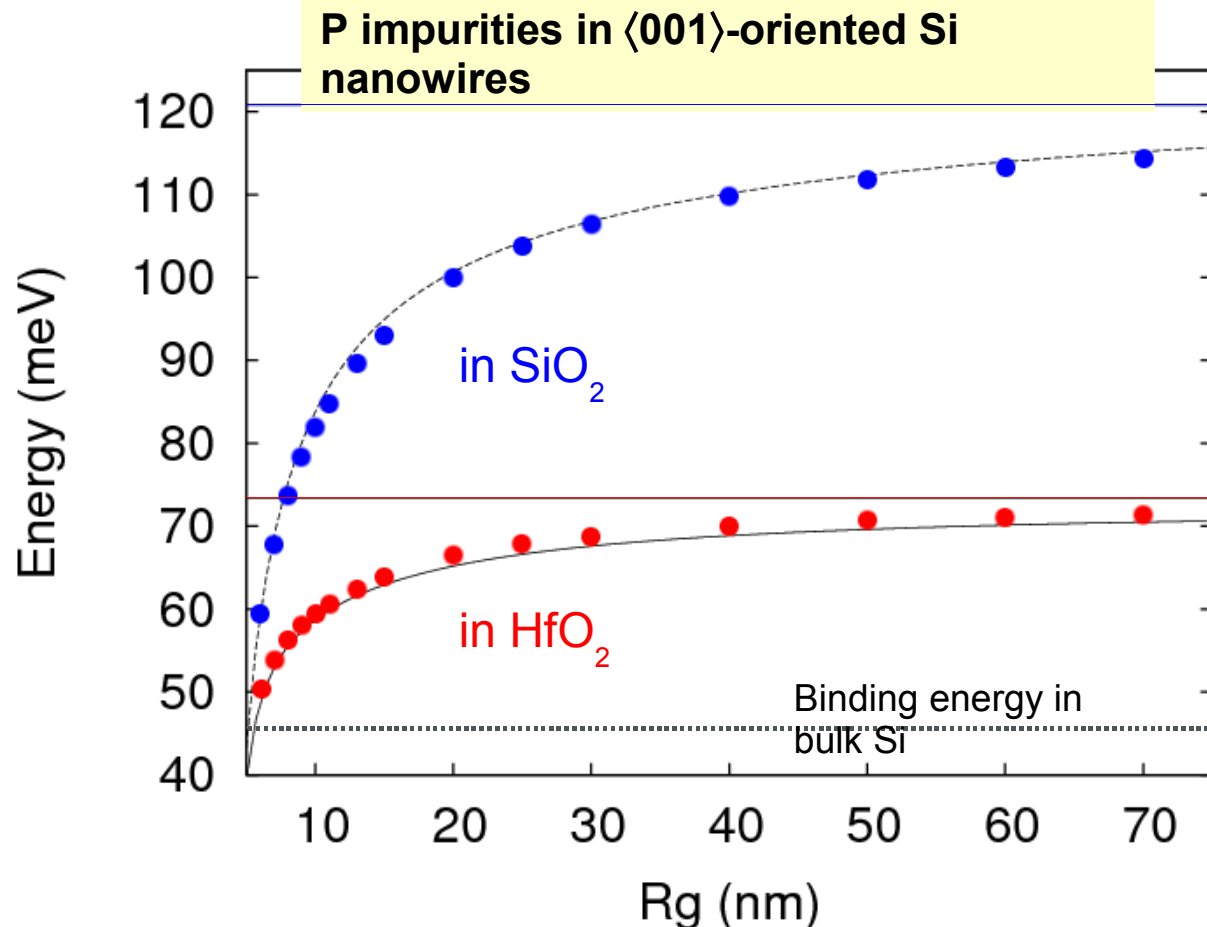
# Si nanowires embedded in $\text{HfO}_2$

- Strong increase of the doping efficiency ( $T = 300 \text{ K}$ ):
  - $P_{\text{ionization}} = 6\%$  @  $R = 5 \text{ nm}$  in vacuum.
  - $P_{\text{ionization}} = 47\%$  @  $R = 5 \text{ nm}$  in  $\text{SiO}_2$ .
  - $P_{\text{ionization}} = 78\%$  @  $R = 5 \text{ nm}$  in  $\text{HfO}_2$ .

## P impurities in $\langle 001 \rangle$ -oriented Si nanowires



# Effect of an « all-around » metallic gate



**R = 5 nm**

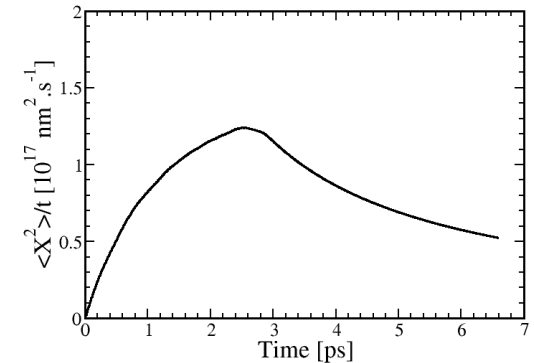
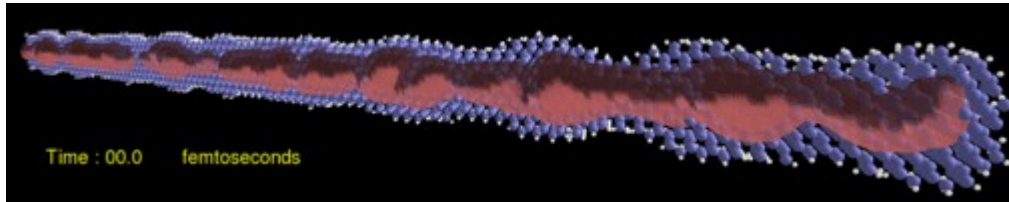
- Strong decrease of the binding energy.
- **Analytical model available for any  $\kappa_{in}$ ,  $\kappa_{out}$ , R and  $R_g \Rightarrow$  Can easily be taken into account in device simulation.**

## Part III :

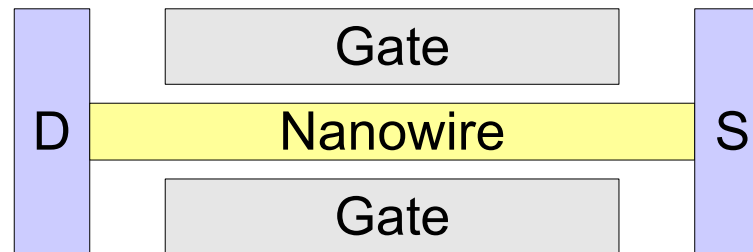
# Transport in silicon nanowires

# The Kubo and Landauer-Büttiker methods

- **Kubo method** : propagate random wavepackets along the nanowires.
  - Yields the « **intrinsic** » transport properties of infinite, disordered nanowires (e.g., mean free paths and mobilities).



- **Green functions method** :
  - Yield the transmission/conductance through a nanowire connected to drain and source **electrodes** (transistor configuration).



- The two methods are complementary and well suited to localized basis sets.



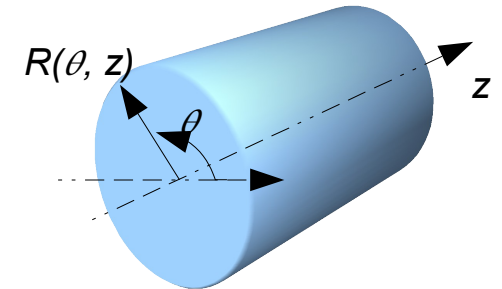
# Application : Surface roughness

- **Disorder** : **Random fluctuations of the radius** of the nanowire, characterized by the auto-correlation function :

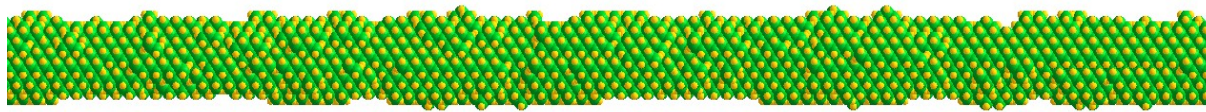
$$\langle \delta R(z, \theta) \delta R(z + \delta z, \theta + \delta \theta) \rangle \equiv \delta R_0^2 e^{-\sqrt{\delta z^2 + R_0^2 \delta \theta^2} / L_r}$$

## Parameters :

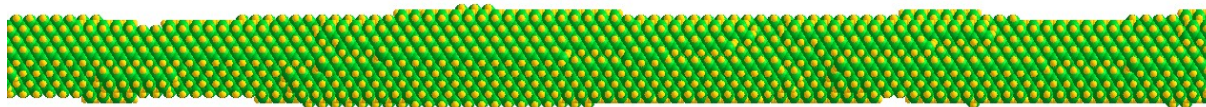
- $R_0$  : average radius.
- $\delta R_0$  : rms fluctuations of the radius.
- $L_r$  : correlation length ( $\sim$  typical size) of the fluctuations.



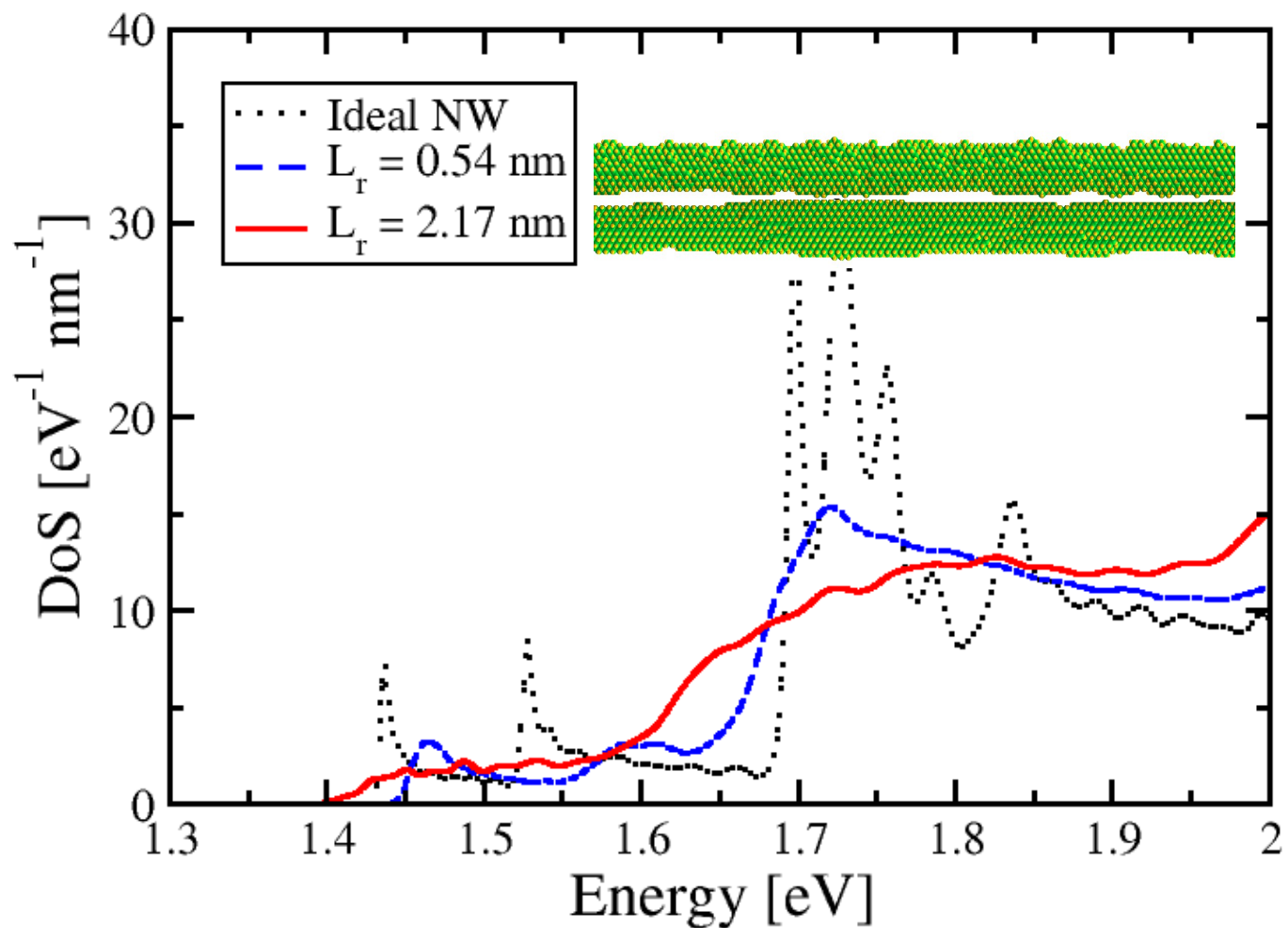
**Si  $\langle 110 \rangle$  nanowires  $R_0 = 1$  nm,  $\delta R_0 = 1$  Å**



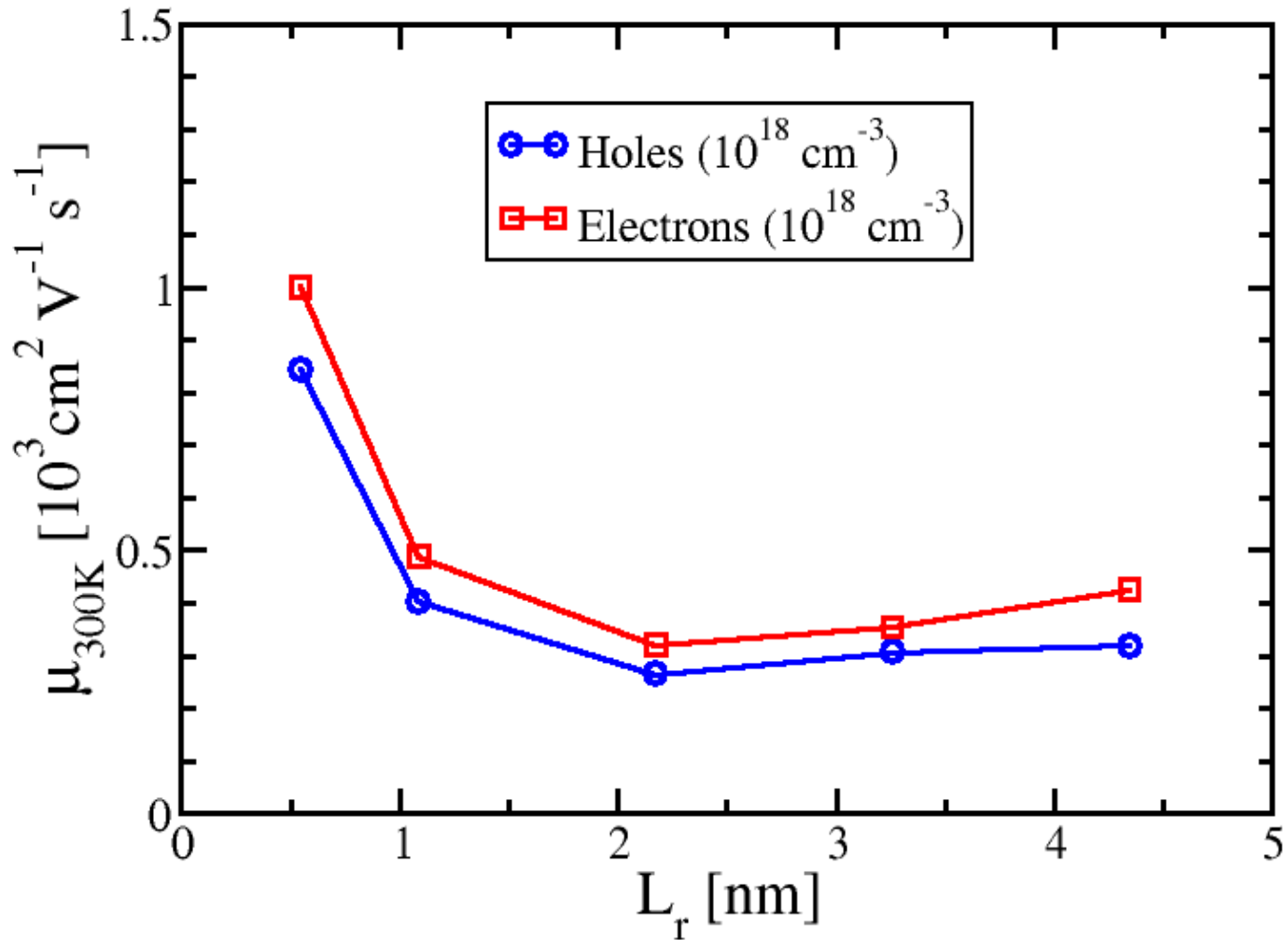
$L_r = 5.4$  Å



$L_r = 21.7$  Å

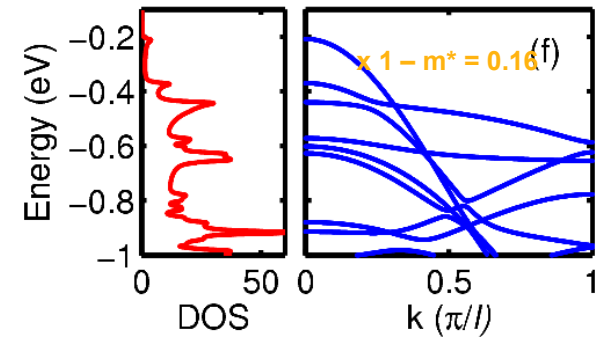
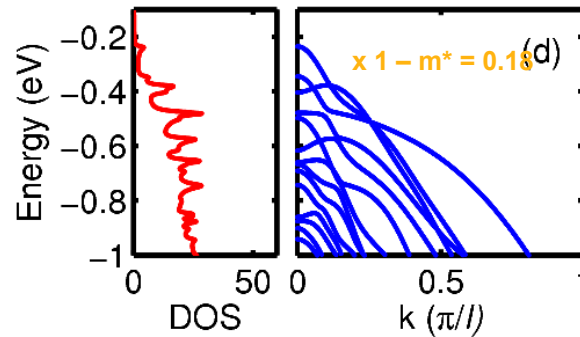
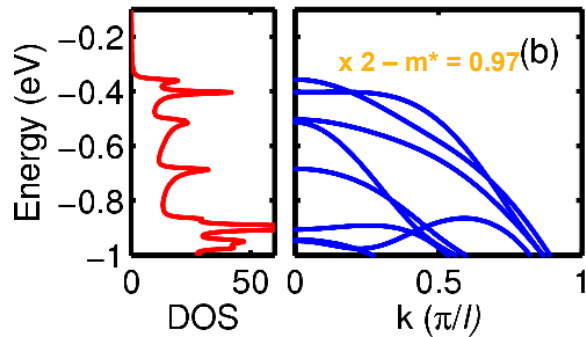
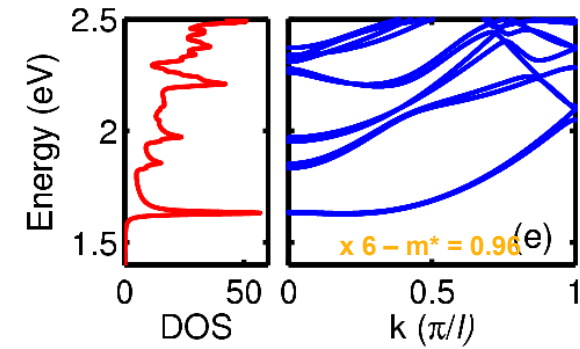
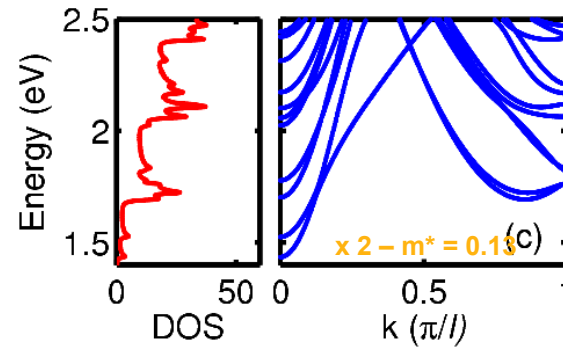
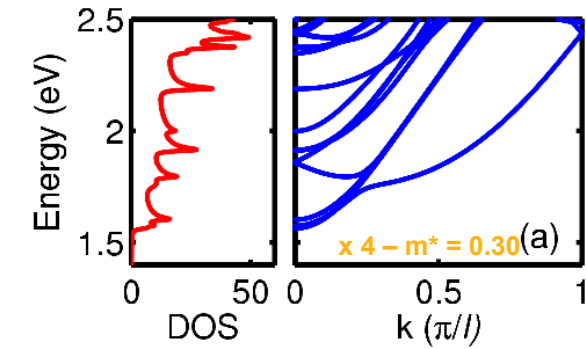
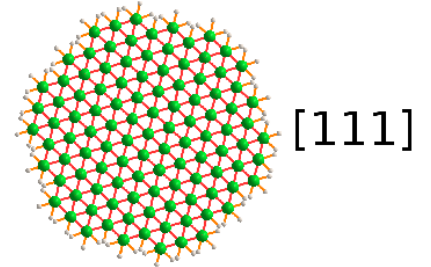
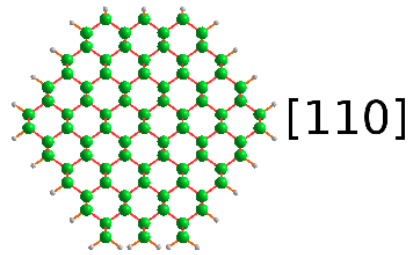
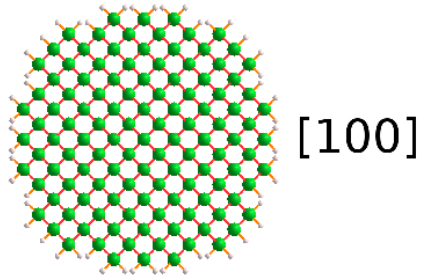


- The Van Hove singularities are smoothed as  $L_r$  is increased.



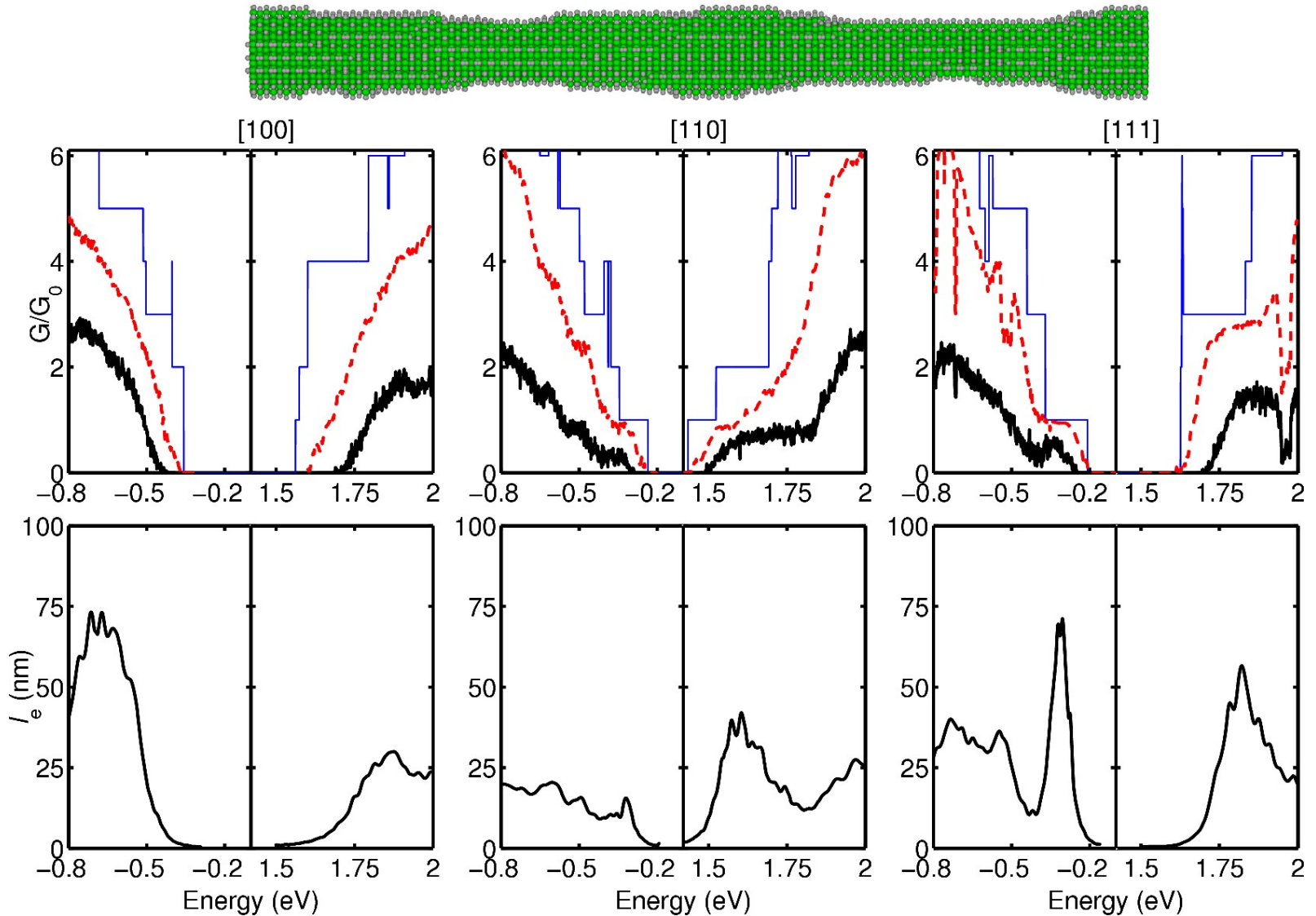
- The mobility of the electrons and holes shows a minimum around  $L_r = 2.5$  nm.

# Band structure of silicon nanowires

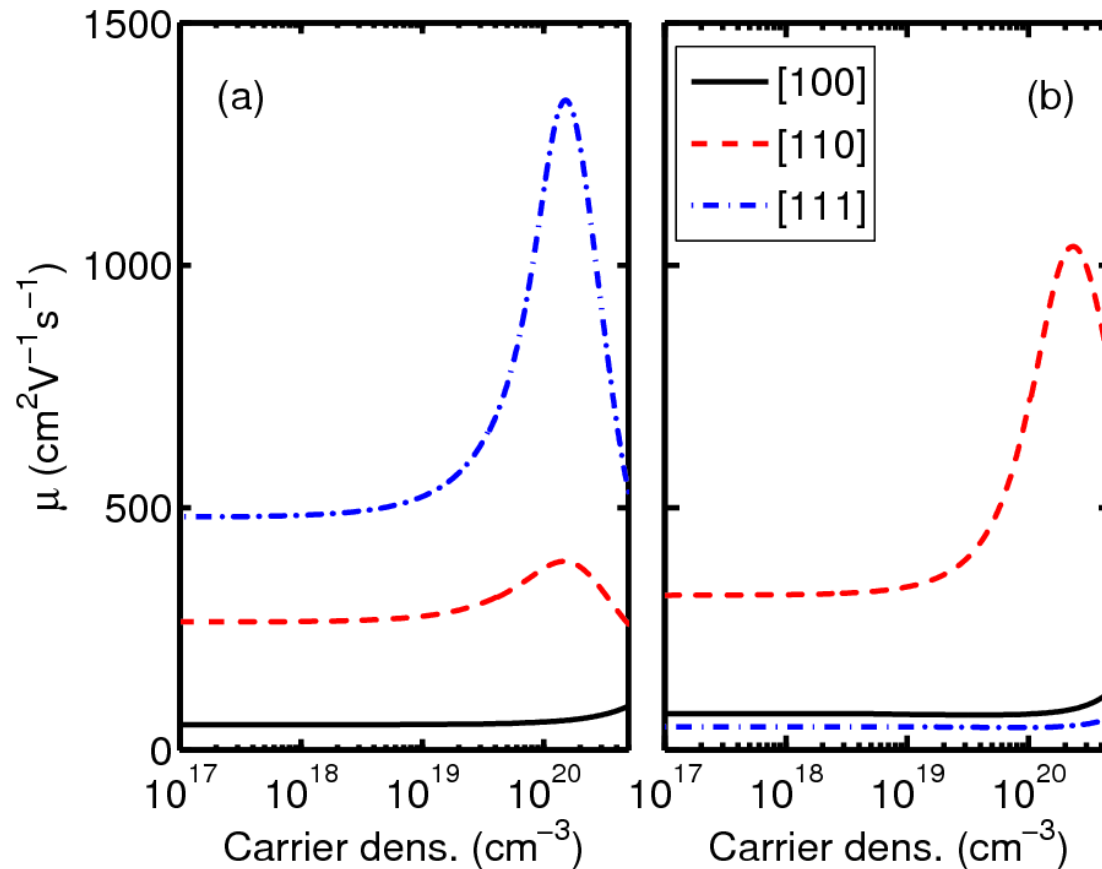


- The band structure of thin Si NWs is strongly dependent on their orientation :
  - Conduction band **valley degeneracy** completely **lifted in [110] Si NWs**.
  - **Lightest hole mass** and largest valence subband splittings **in [111] Si NWs**.

# Transport properties of [100], [110] and [111] oriented Si NWs



# Mobility as a function of Si NW orientation



- In agreement with the trends evidenced on the band structures,
  - [111] is the best orientation for hole transport.
  - [110] is the best orientation for electron transport.



- Strain relaxation is very **efficient** in nanowire heterostructures.
- In short barriers and quantum-dots **the potential barrier is lowered**.
- **The binding energy of donor and acceptor impurities can be greatly enhanced in small semiconductor nanowires, which decreases the doping efficiency.**
- **The binding energy of the impurities however depends on the dielectric environment of the nanowires** (through screening and polaronic effects). High- $\kappa$  dielectrics and metallic gates can help to increase the doping efficiency.
- **The transport properties of thin silicon nanowires strongly depend on their orientation :**
  - The best orientation for **electron** transport is **[110]**.
  - The best orientation for **hole** transport is **[111]**.

This can be related to the anisotropy of the band structure of silicon and should not be much sensitive to the nature of the scattering mechanisms.