

## Explicit computation of Coulomb and exchange interactions for N-electrons in open quantum systems using Bohm trajectories

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From a computational point of view, the direct solution of the many-particle Schrödinger equation is inaccessible for more than very few electrons. This issue is at the heart of almost all the unsolved problems in quantum transport. Recently, a novel many-particle quantum transport formalism using Bohm trajectories has been presented for dealing with Coulomb and exchange interaction among electrons [1]. We discuss the computational burden associated with the explicit consideration of the electron spin in the previous formalism [1]. In particular, we provide a numerical justification that shows the viability of the previous formalism for a system with a large ( $N \approx 100$ ) number of electrons. We consider a system of  $N$  electrons described by a many-particle wave-function  $\Psi(\vec{r}_1, \vec{r}_2, \vec{r}_3 \dots \uparrow_1, \downarrow_2, \downarrow_3 \dots; t)$  with  $\vec{r}_i$  the electron position and  $\uparrow_i / \downarrow_i$  its (up/down) spin. We use an uncoupled spin-base which is adequate for (non-conservative spin) open systems. In the previous formalism [1], the Bohm velocity of each electron has to be computed directly from the many-particle wave-function. Since this formalism is intrinsically time-dependent [1], the standard simplifications in the computation of the many-particle wave-function, obtained from Hamiltonian (orthogonal) eigenstates, cannot be used. In addition, the explicit evaluation of  $N!N!$  products of permutations for the computation of the many-particle system is intractable for more than very few electrons because of computational limitations (note that  $8!^2 = 40320^2$ ). The previous computational limitation is overcome by computing the many-particle velocity with the assumption that the many-particle wave-function can be separated into a product of spin-up ( $\uparrow$ ) and spin-down ( $\downarrow$ ) many-particle wave functions:

$$\Psi(\vec{r}_1, \vec{r}_2, \vec{r}_3 \dots \uparrow_1, \downarrow_2, \downarrow_3 \dots, t) \approx \Psi(\vec{r}_1, \vec{r}_4 \dots \uparrow_1, \uparrow_4 \dots, t) \cdot \Psi(\vec{r}_2, \vec{r}_3 \dots \downarrow_2, \downarrow_3 \dots, t). \quad (1)$$

Then, the numerical difficulties in the computation of the many-particle Bohm velocity disappear because it can be computed from a complex matrix (Slater) determinant. In order to numerically verify the correctness of our assumption, we compute the Bohm velocity associated to electron 1 in four different (exchange-interacting) situations. We have defined the parameter  $d$  as a normalized (i.e. without units) phase-space distance [2] between electron 1 and the others (see insets in all Figures). We have chosen arbitrary initial Gaussian wave-packets. In Fig. 1, we show the Bohm velocity (with an approximate value of  $6 \times 10^4$  m/s) for one independent (spin-up) electron. In Fig. 2, we plot the same Bohm velocity for the same electron when other 4 exchange-interacting electrons are present. As we decrease the distance  $d$  among electrons, the Bohm velocity becomes very different from Fig. 1 as a consequence of the Pauli (Exclusion) Principle. In Fig. 3, we show one spin-up electron and two spin-down electrons. Surprisingly, when Fig. 3 is compared with Fig. 1, an exchange interaction between spin-up and spin-down electrons is present due to the non-orthogonal spatial behavior of our "arbitrary" wave-functions. However, the effect of such spin-up-spin-down interaction in the Bohm velocity is quite small when compared to its effect with the spin configuration of Fig. 4, where the 3 spin-up electrons of Fig. 2 are considered. The strong resemblance between the Bohm velocities of Figs. 2 and 4 for the different values of  $d$  provides a numerical justification

of expression (1) for the computation of many-particle Bohm velocities. The same result is obtained for many other spin schemes.

In conclusion, we present a powerful proposal to study quantum transport with the explicit consideration of Coulomb and exchange interaction among electrons. We show the practical numerical viability of our proposal for large,  $N \approx 100$ , number of electrons. Since the algorithm deals with time-dependent Schrödinger equations in terms of quantum trajectories, it can be applied for the computation of the average current or its fluctuations [3] in zero or high frequency [4] quantum scenarios.

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**References:**

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**Figures:**

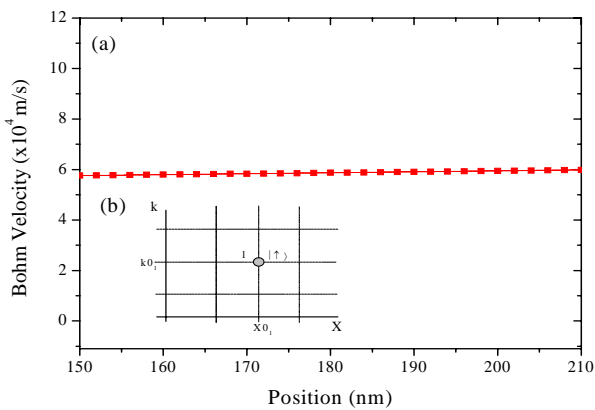


Fig. 1. (a) Bohm velocity for an independent electron. (b) Schematic representation of the system for an electron where we indicate the central value of the X<sub>0</sub> and wave-vector K<sub>0</sub>.

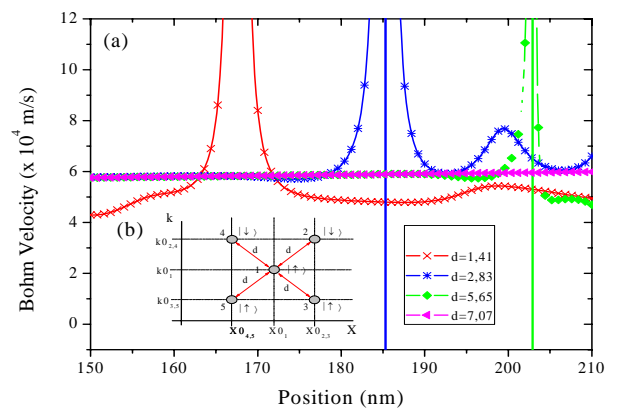


Fig. 2. (a) Bohm velocities for 1-electron using different values of  $d$  for a system of 5 electrons (3 spin-up and 2 spin-down). (b) In this scheme we indicate the central value of the X<sub>0</sub> and wave-vector K<sub>0</sub> of these electrons.

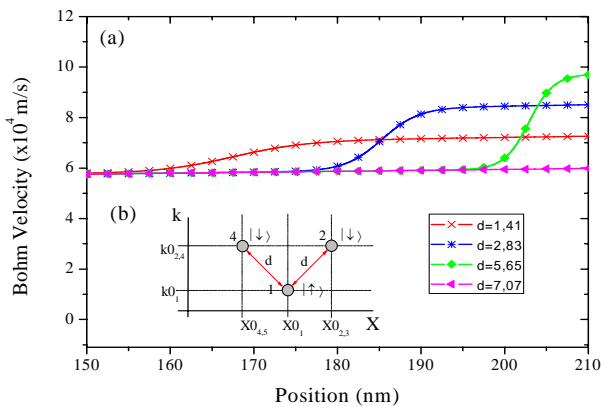


Fig. 3. (a) Bohm velocities for 1-electron using different values of  $d$  for a system of 3 electrons (1 spin-up and 2 spin-down). (b) In this scheme we indicate the central value of the X<sub>0</sub> and wave-vector K<sub>0</sub> of these electrons.

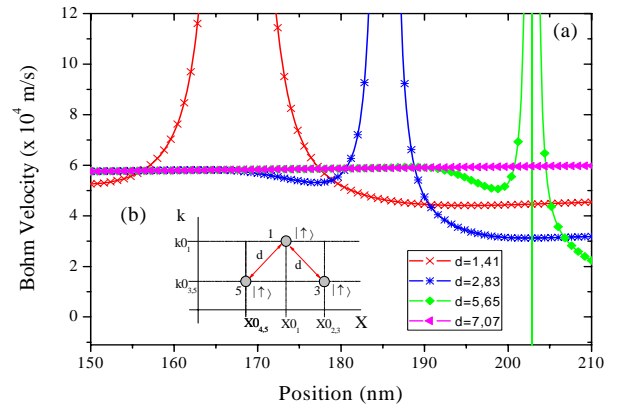


Fig. 4. Bohm velocities for 1-electron using different values of  $d$  for a system of 3 electrons (spin-up). (b) In this scheme we indicate the central value of the X<sub>0</sub> and wave-vector K<sub>0</sub> of these electrons.