



UNIVERSITAT DE BARCELONA



# Exchange bias in core/shell magnetic nanoparticles: experimental results and numerical simulations

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# I. Motivation

- Magnetic nanoparticles
- Why interesting? – Size effects (1940s):
  - enhanced properties with respect to bulk
  - finite size and surface effects + interactions (dipolar)

Feature size  $\leq$  correlation length

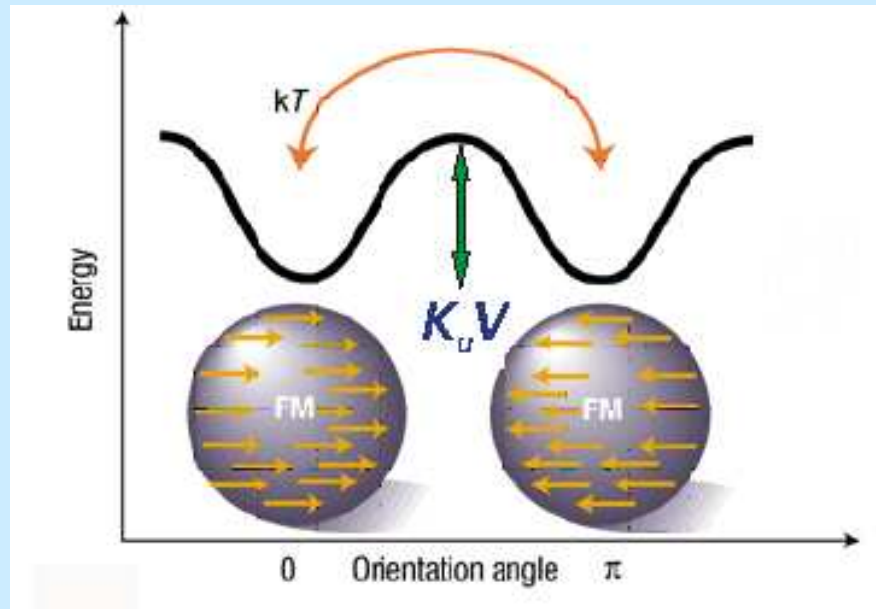
Exchange length, domain wall width and single domain radius

- Why useful? – Technological applications:
  - Magnetic recording
  - Biomedical applications

# Magnetic recording vs biomedical applications

The superparamagnetic limit: thermal fluctuations overcome the magnetic anisotropy barrier

$$t = t_0 e^{K_u V / kT}$$



Magnetic recording ( $t \approx 10$  years)

With the increase in recording density (reduction in volume), we need to keep the magnetic stability (patterned media, high anisotropy materials, exchange bias...)

Requirements for bio-applications (typically)

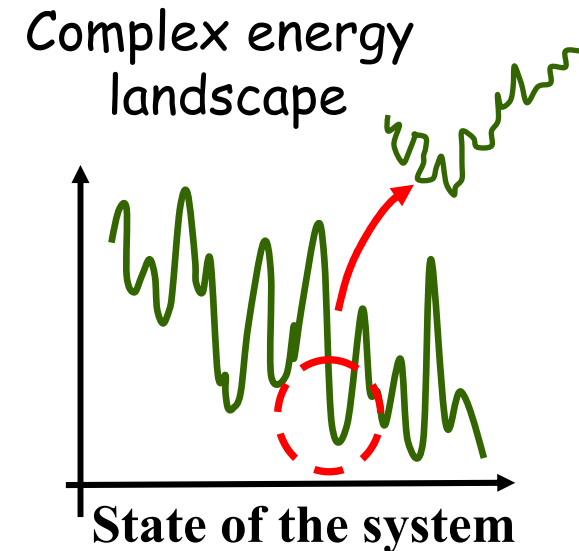
Superparamagnetic behavior

High magnetization

Limiting size (*in vivo*)

Biocompatibility and functionality

## II. Magnetic nanoparticles: Finite-size, surface effects and collective behavior



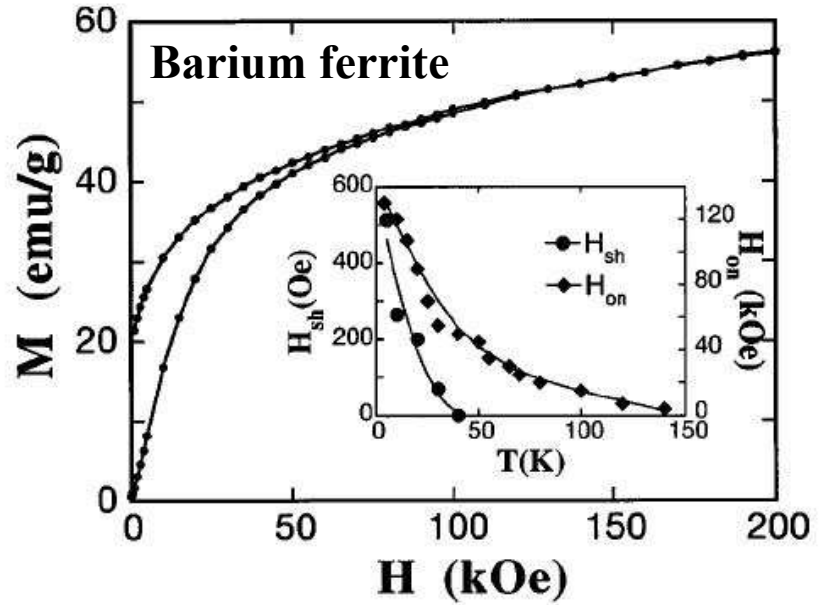
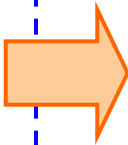
### Review Papers

Finite size effects in fine particles: magnetic and transport properties  
X. Batlle and A. Labarta,  
*Journal of Physics D: Applied Physics* **35**, R15 (2002)

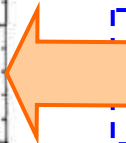
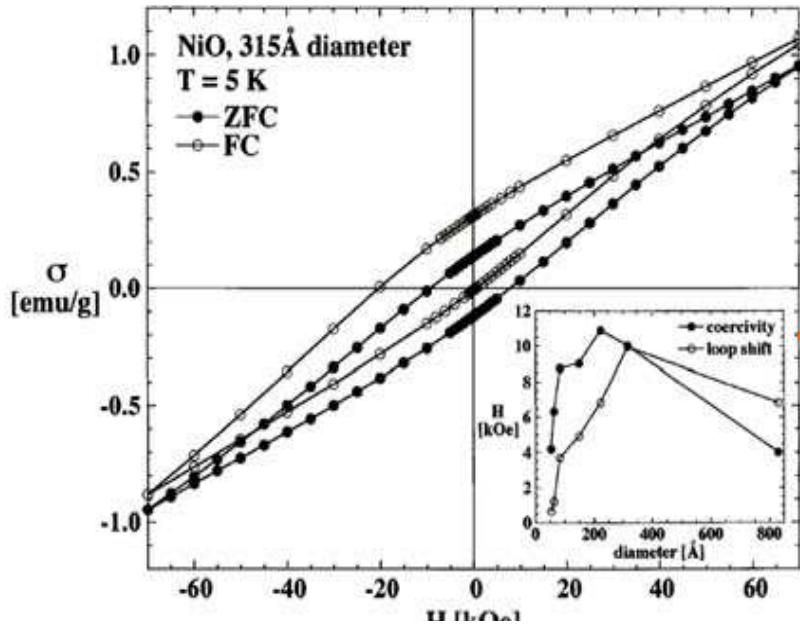
From finite-size and surface effects to glassy behaviour in ferrimagnetic nanoparticles

O. Iglesias, X. Batlle and A. Labarta  
*Surface effects in magnetic nanoparticles*; D. Fiorani Editor; Springer (2006)

Low saturation magnetization,  
high field irreversibility,  
high closure fields  
(broken symmetry at surface)



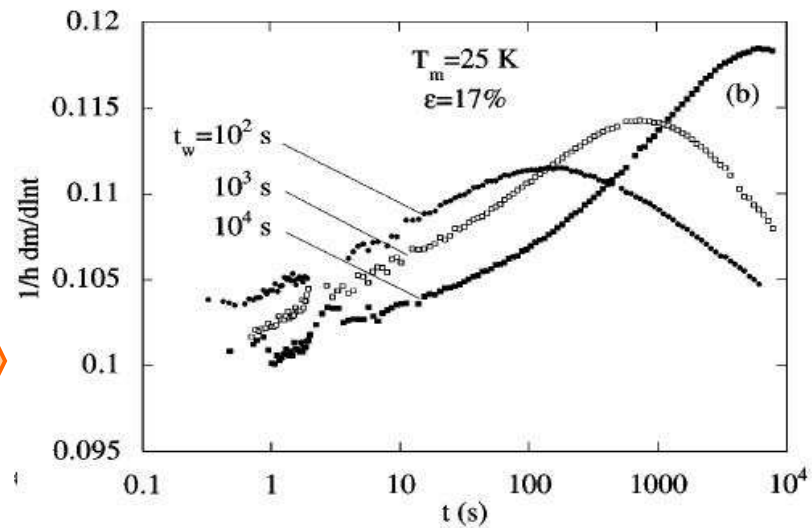
M. García et al., PRB 59, 13594 (99) UB



Shifted hysteresis loops after FC

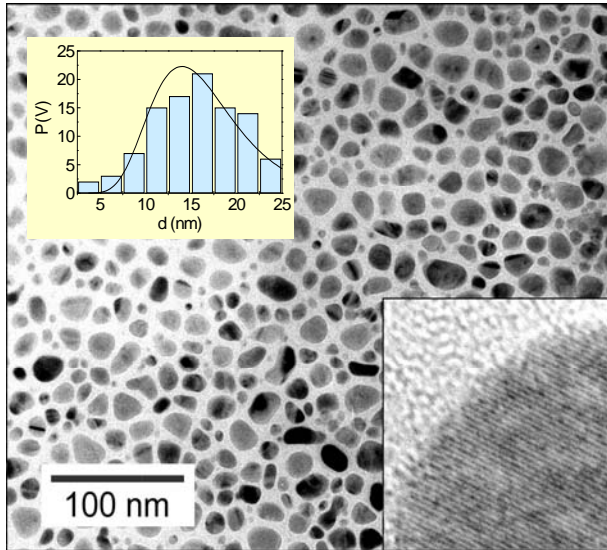
R. H. Kodama et al., PRL 79, 1393 (1997)

Ageing effects  
(magnetic properties  
depend on waiting time)



T. Jonsson et al. PRL 75, 4138 (95)

# Size distribution



**Co-ZrO<sub>2</sub>**

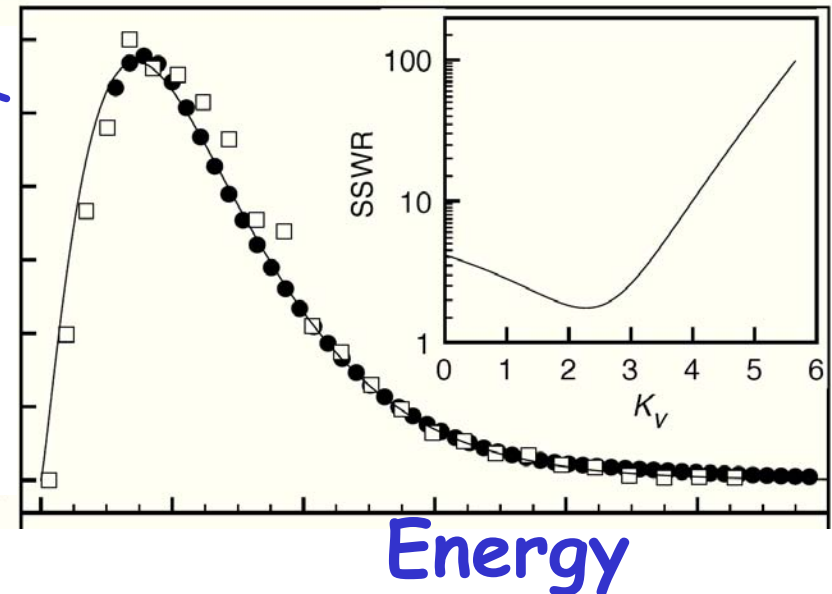
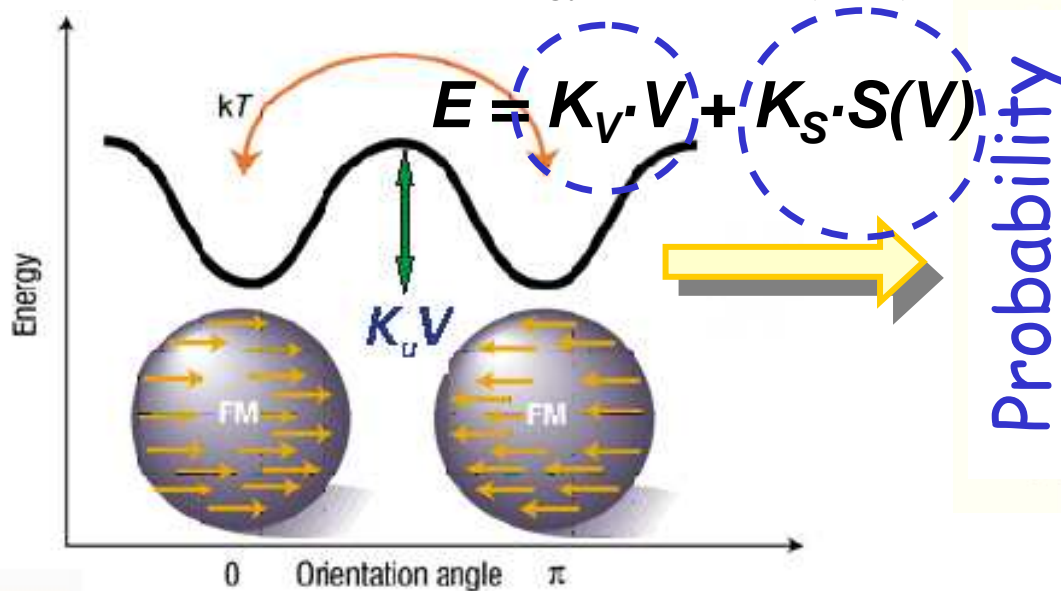
B.J. Hattink et al., PRB **67**, 033402 (2003) **UB**

N. Pérez et al., Nanotechnology **19**, 475704 (2008) **UB**

Non-uniform distribution  
(log-normal)

$$f(V) = \frac{1}{\sqrt{2\pi}\sigma V} \exp\left(-\frac{\ln^2(V/V_0)}{2\sigma^2}\right)$$

Effective distribution of  
energy barriers for  
magnetization reversal

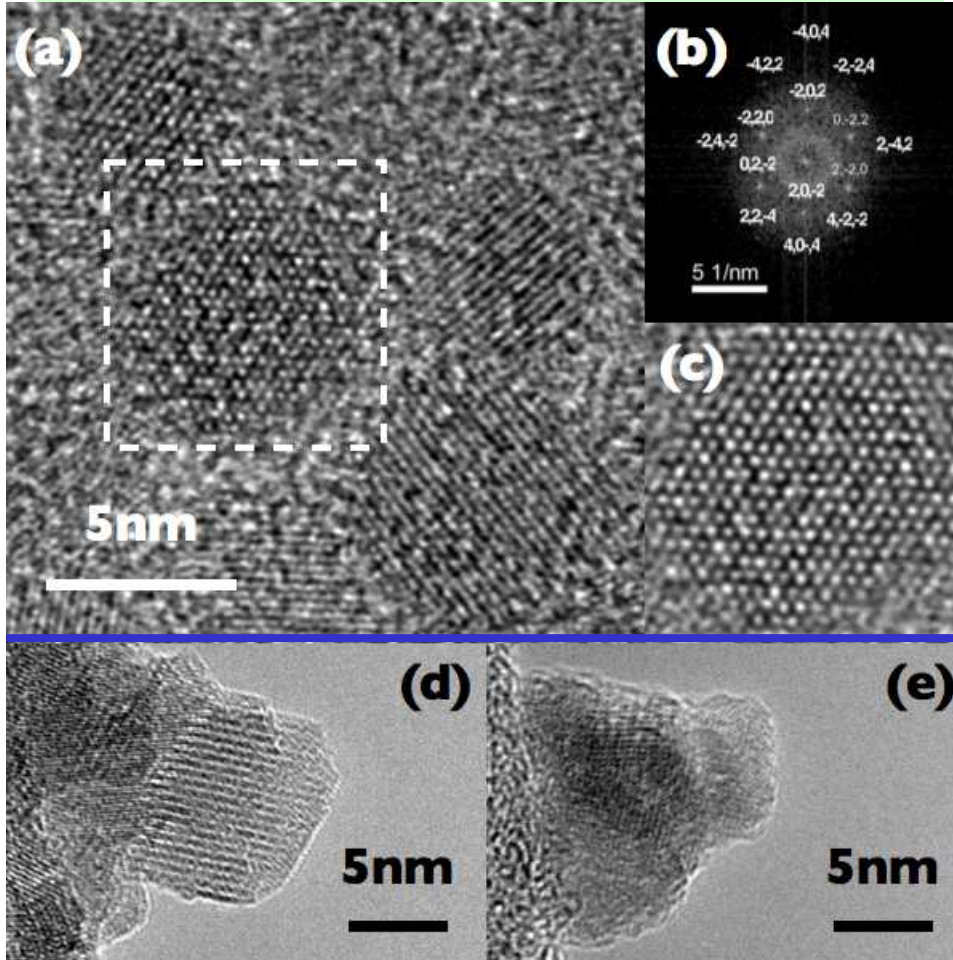




# COATING EFFECTS

5 nm  $\text{Fe}_3\text{O}_4$  particles

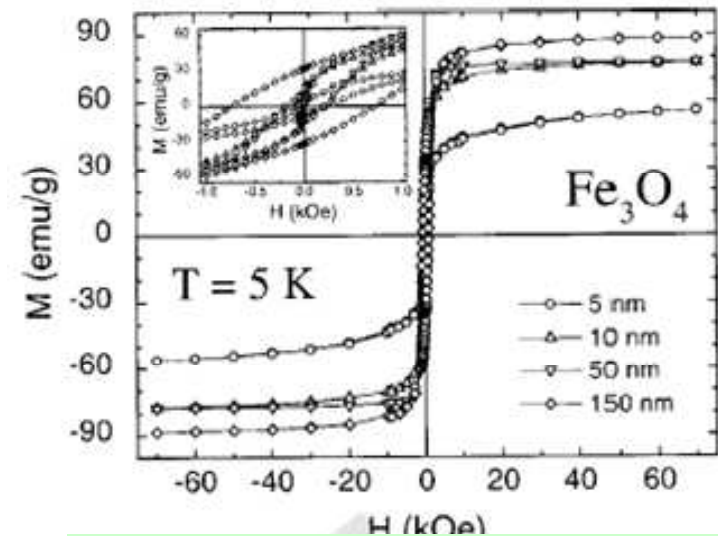
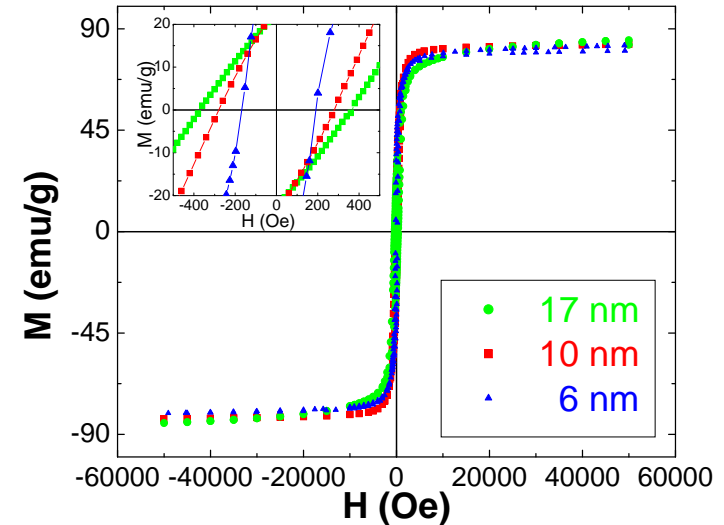
Covalent bond (oleic acid)



Adsorbed coating (PVA)

N. Pérez et al., APL 94, 093108 (2009) UB

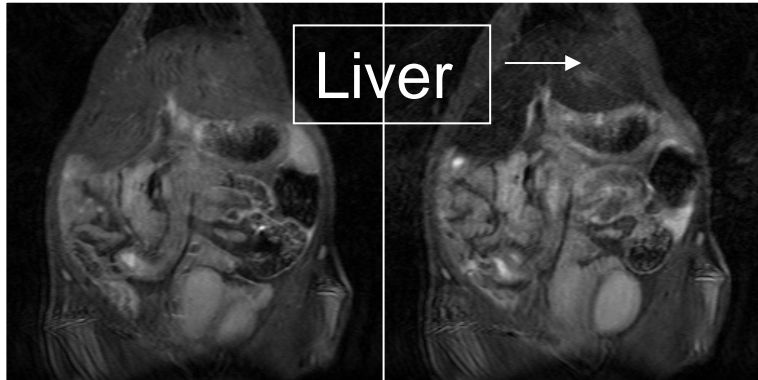
Size-independent,  
bulk magnetization



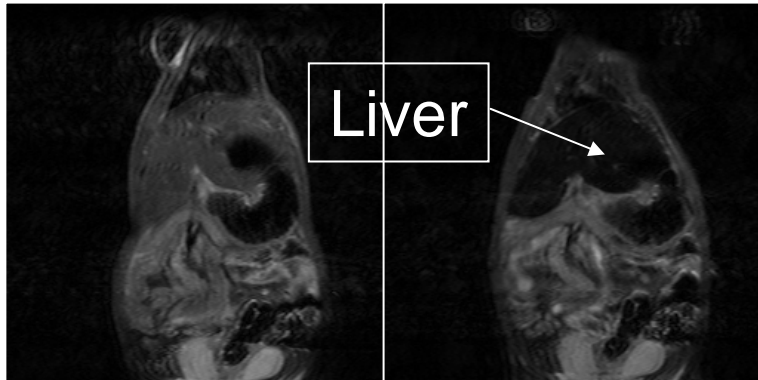
Size-dependent,  
reduced magnetization

# MRI

BEFORE AFTER (9 nm)



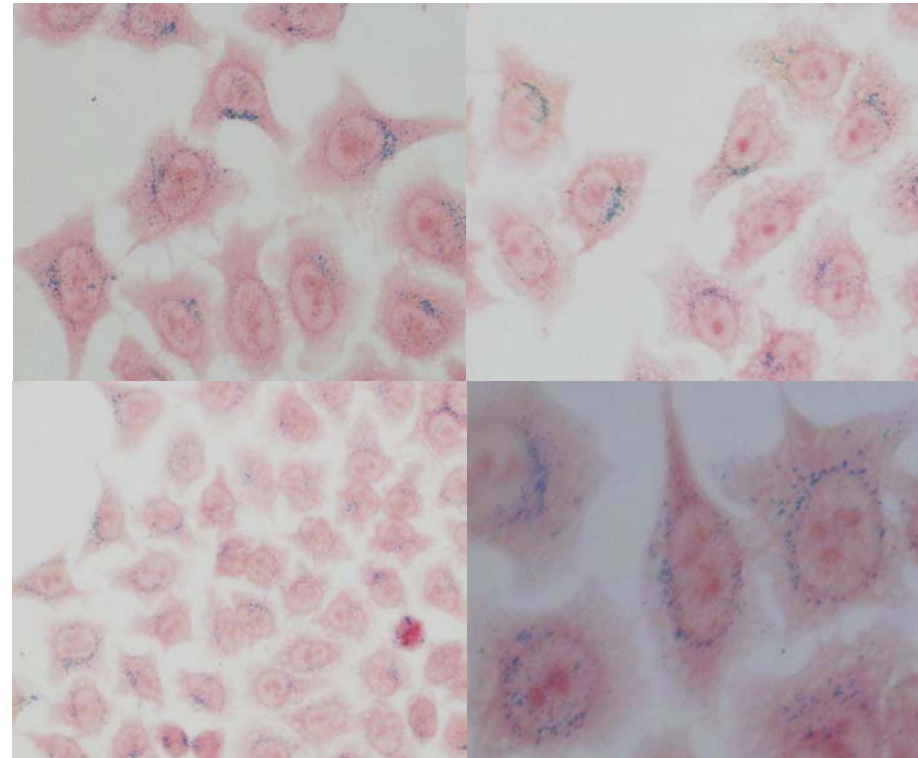
BEFORE AFTER (4 nm)



**Contrast enhancement**

M.P. Morales et al., (submitted) **UB**

# Incubation with HeLa cells



**Internalization around the cell membrane**



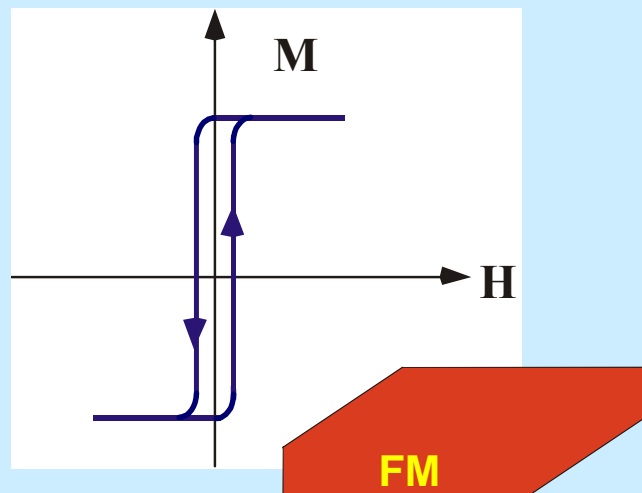
# III. Exchange bias: A (very short) survey

Review Paper

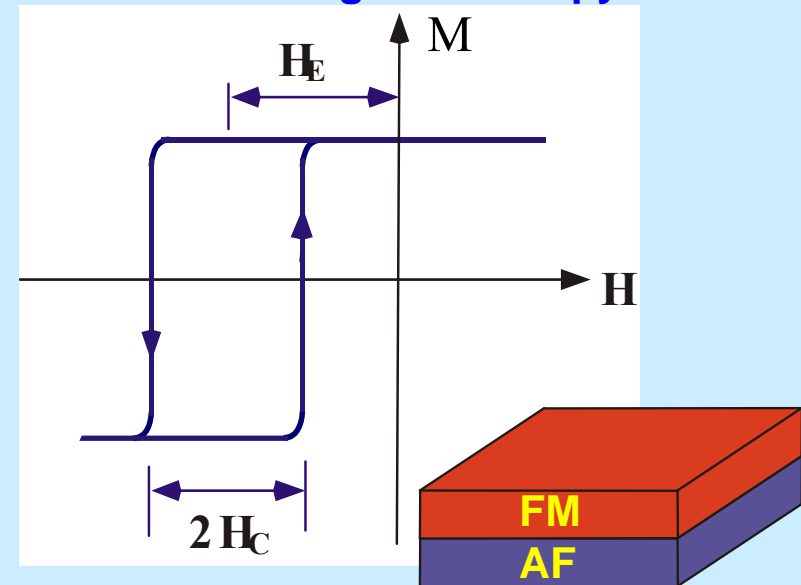
Exchange bias phenomenology and models of core/shell nanoparticles  
O. Iglesias, A. Labarta and X. Batlle  
Journal of Nanosciences and Nanotechnology **8**, 2761 (2008)

# Reminder

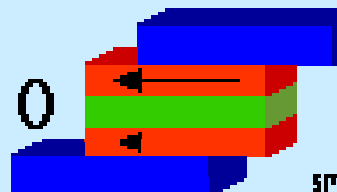
- Exchange bias: unidirectional anisotropy induced by the AF into the FM via exchange coupling at the interface



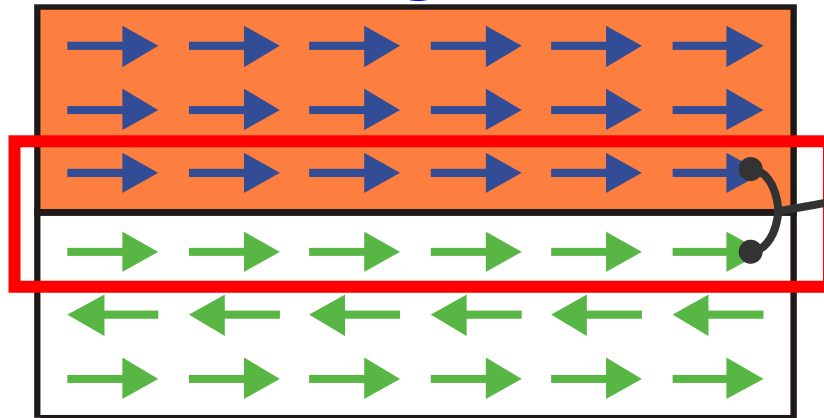
$M(H)$  after field cooling  $T_N < T < T_C$   
AFM with high anisotropy



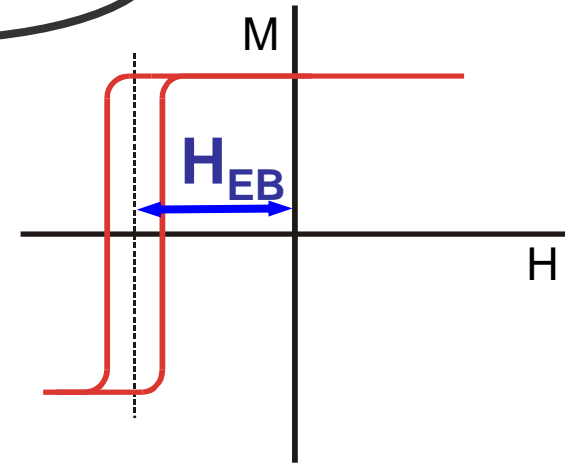
- Reference spin state (read heads, MRAM and magnetic sensors)



# Microscopic origin of exchange bias

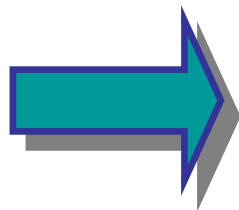


exchange coupling



$$H_{EB} = n \frac{JS_{AF}S_{FM}}{t_{FM}M_{FM}}$$

$n$  ?



1. Uncompensated (3 – 7%) pinned spins in the AF, either at the interface or far from it.

Ohldag, PRL **91**, 017203 (2003)

Kappenberger, PRL **91**, 267202 (2003)

Roy, PRL **95**, 047201 (2005) **UB**

2. Asymmetric reversal due to broken symmetry, leading to *in-volume* domains in the FM.

Li, PRL **96**, 217205 (2003) **UB**

Morales, APL **89**, 072504 (2006); **95** (2009) **UB**

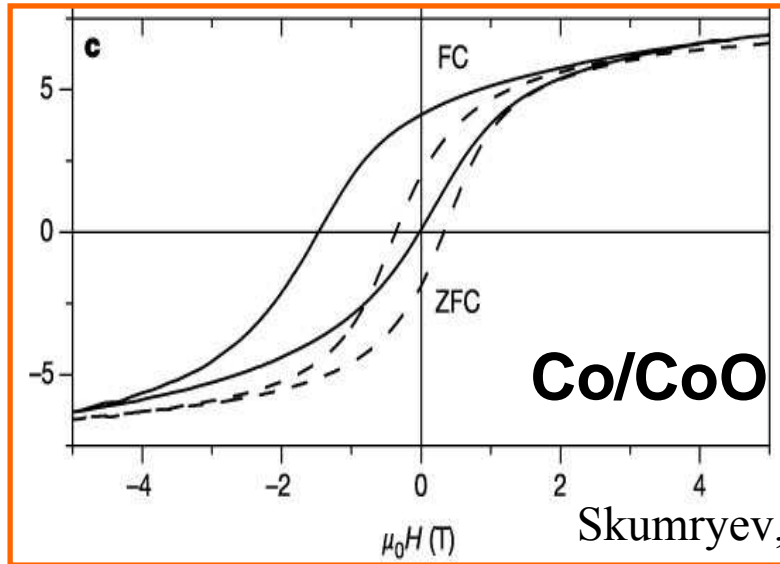
3. Key role of relative size of *in-plane* FM/AF domains

Roshchin, EPL **71**, 297 (2005) **UB**

Petracic, APL **87**, 222509 (2005) **UB**

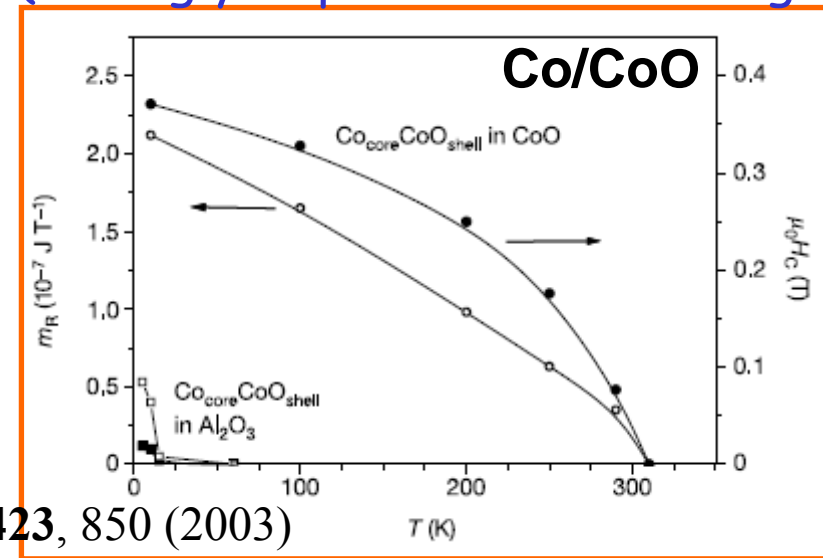
# Phenomenology in core/shell NP

Shifted loops,  
increased  $H_c$  (spin dragging)

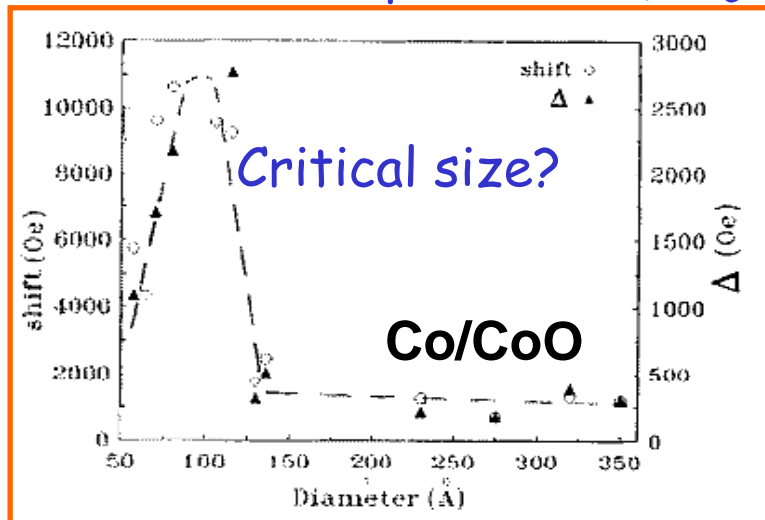


Skumryev, Nature **423**, 850 (2003)

Increased blocking temperature  
(strongly dependent on coverage)

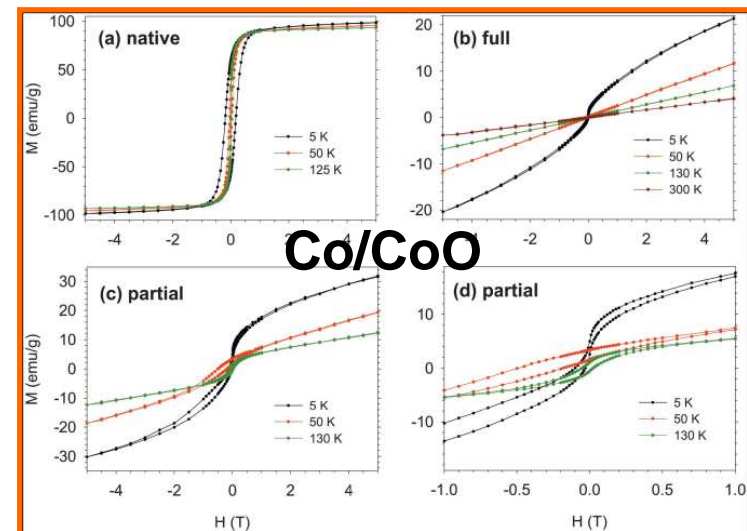


Particle size dependence ( $\sim D_c^{-1}$ )



Gangopadhy, JAP **73**, 6964 (1993)

Oxidation state (FM-AF ratio)



Tracy, PRB **72**, 064404 (2005)

## IV. Results:

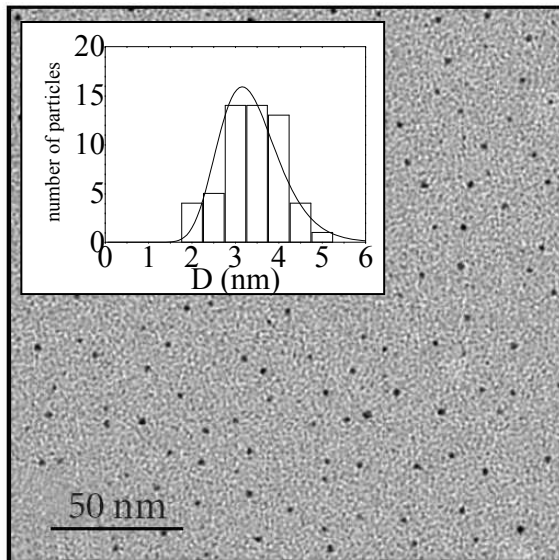
# EB in Co-CoO nanoparticles embedded in a matrix

M. Kovylna, M. García del Muro, Z. Konstantinović, O. Iglesias  
M. Varela, A. Labarta and X. Batlle,

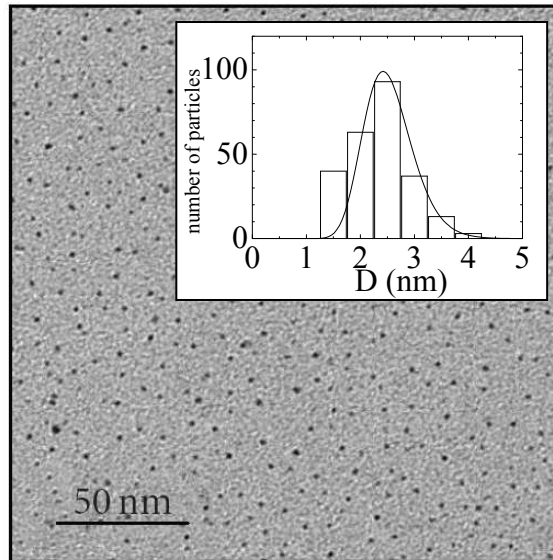
Nanotechnology **20**, 175702 (2009)



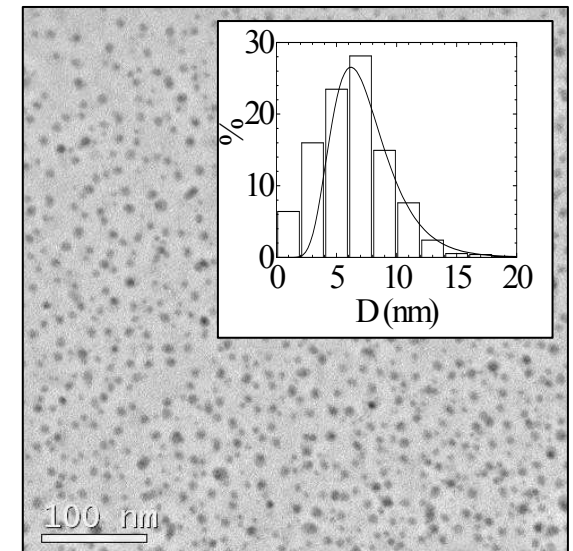
# $\text{Co}_x\text{-(ZrO}_2\text{)}_{1-x}$ thin films by pulsed laser ablation



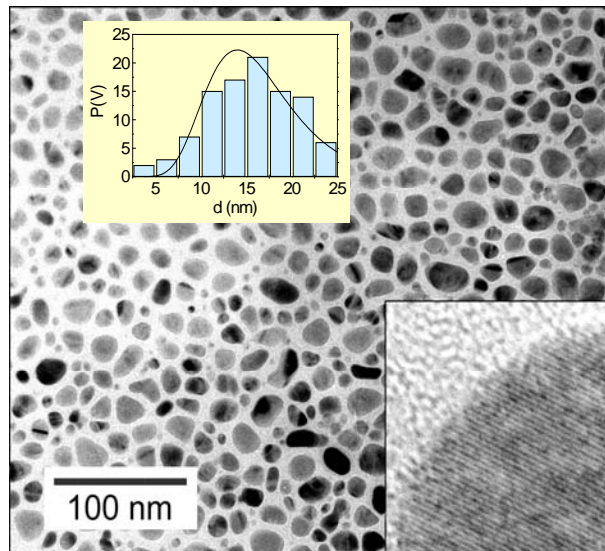
$x$  (% vol.) = 0.06



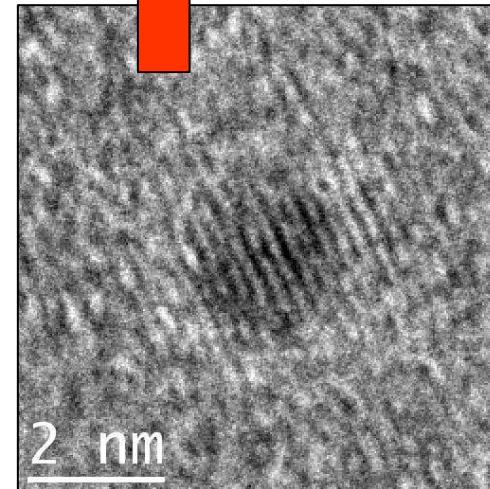
$x$  (% vol.) = 0.12



$x$  (% vol.) = 0.25



$x$  (% vol.) = 0.30

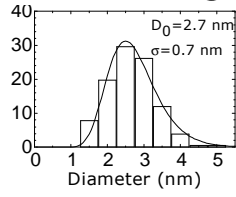
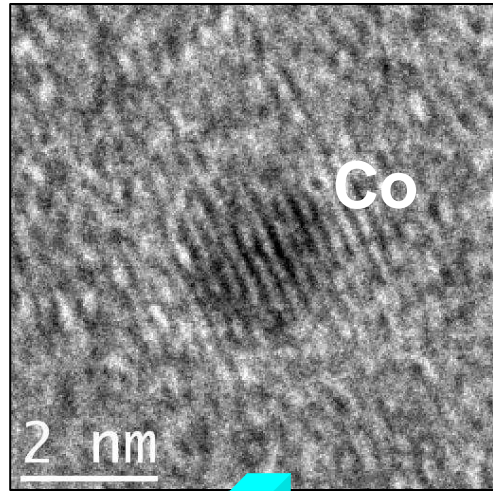


Z. Konstantinovic et al., Nanotechnology **17**, 4106 (2006) **UB**

Z. Konstantinovic, APL **91** 052108 (2009); **90**, 182506 (2009) **UB**

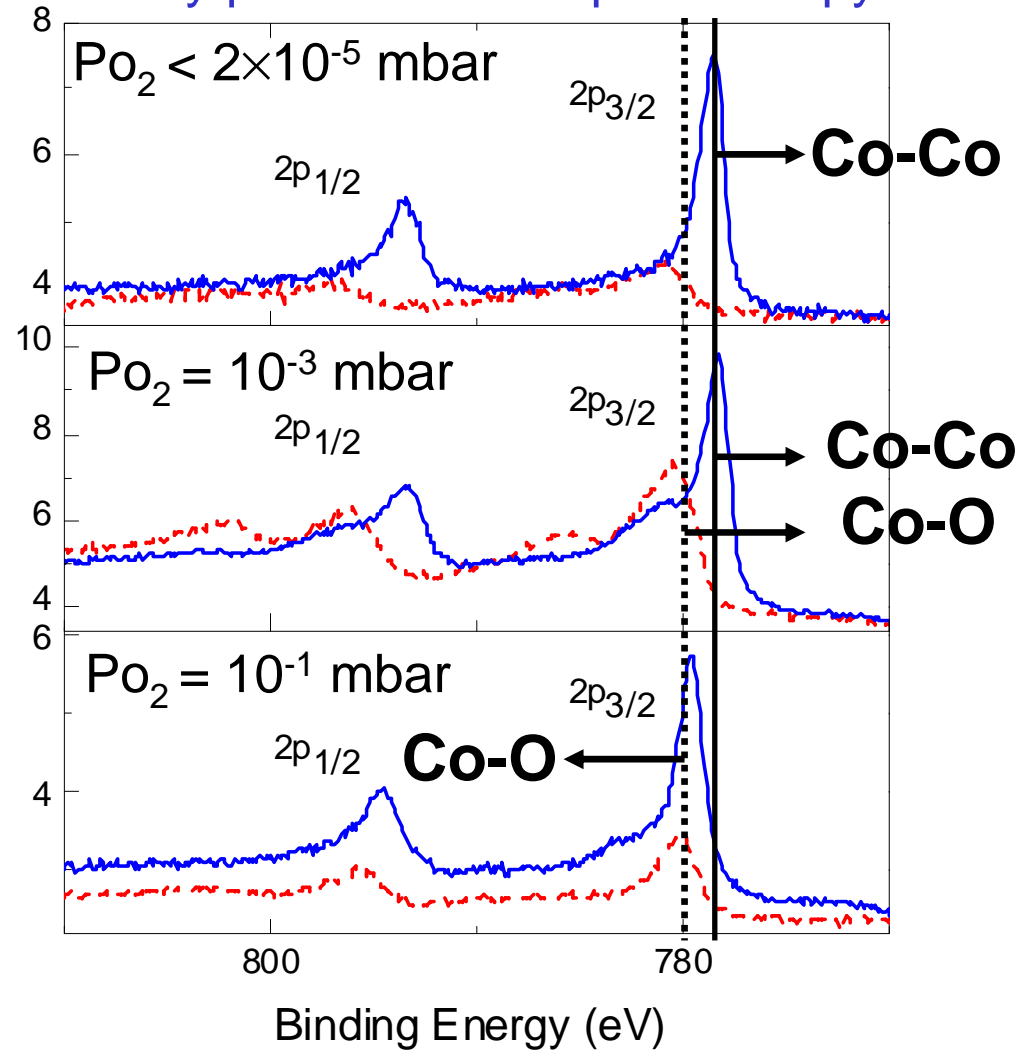
B.J. Hattink et al., PRB **73**, 45418 (2006); **79**, 94201 (2009) **UB**

$P_{O_2} < 2 \times 10^{-5}$  mbar (base)

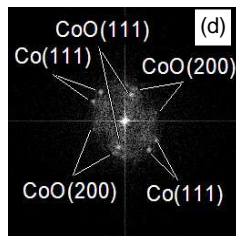
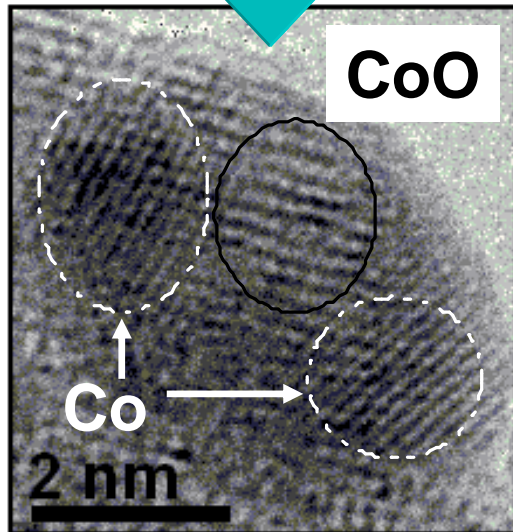


Co-ZrO<sub>2</sub> thin films deposited under oxygen pressure ( $x_v(\text{Co}) \approx 0.2$ )

X-ray photoelectron spectroscopy

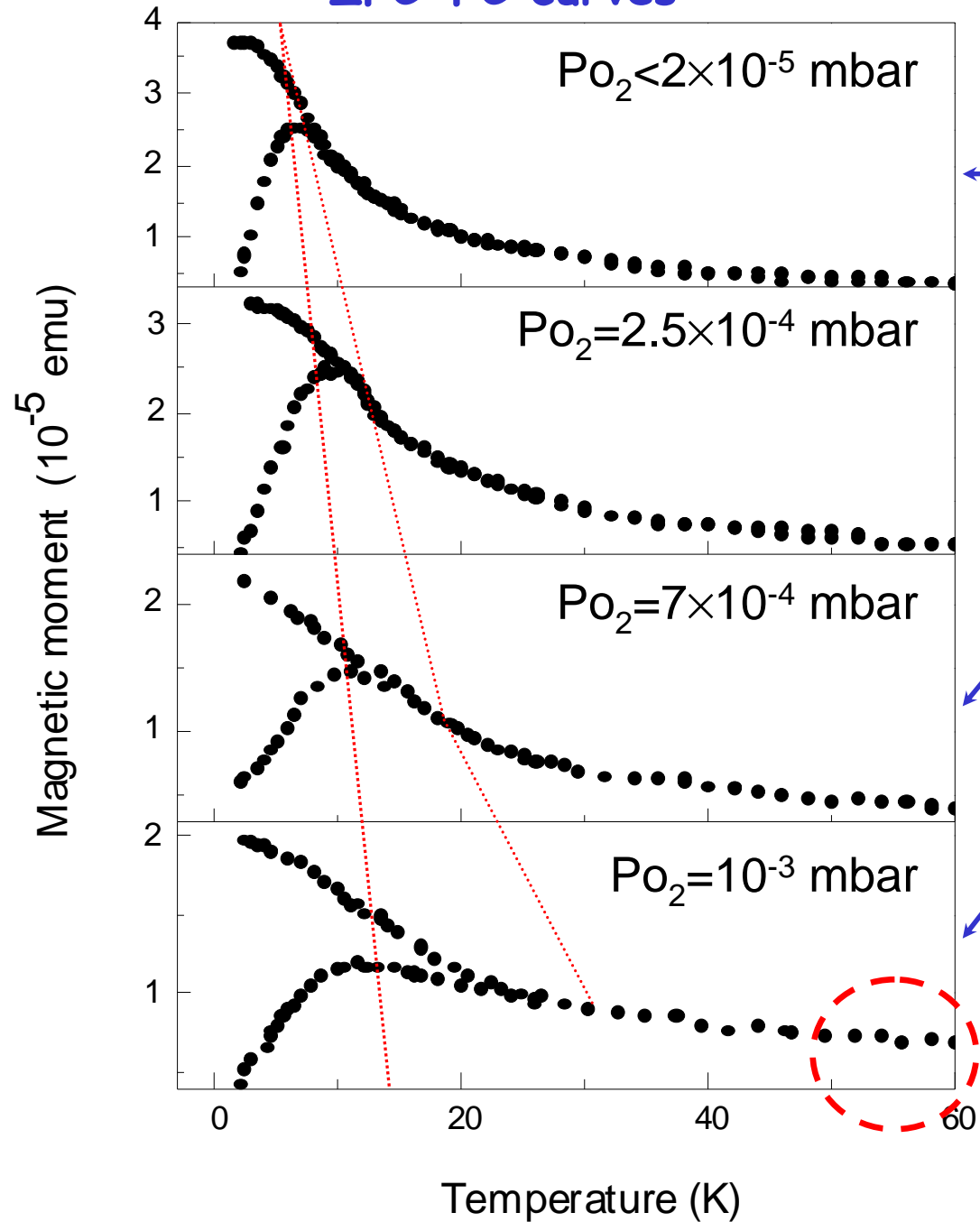


Increasing oxygen pressure  
(partial coverage)

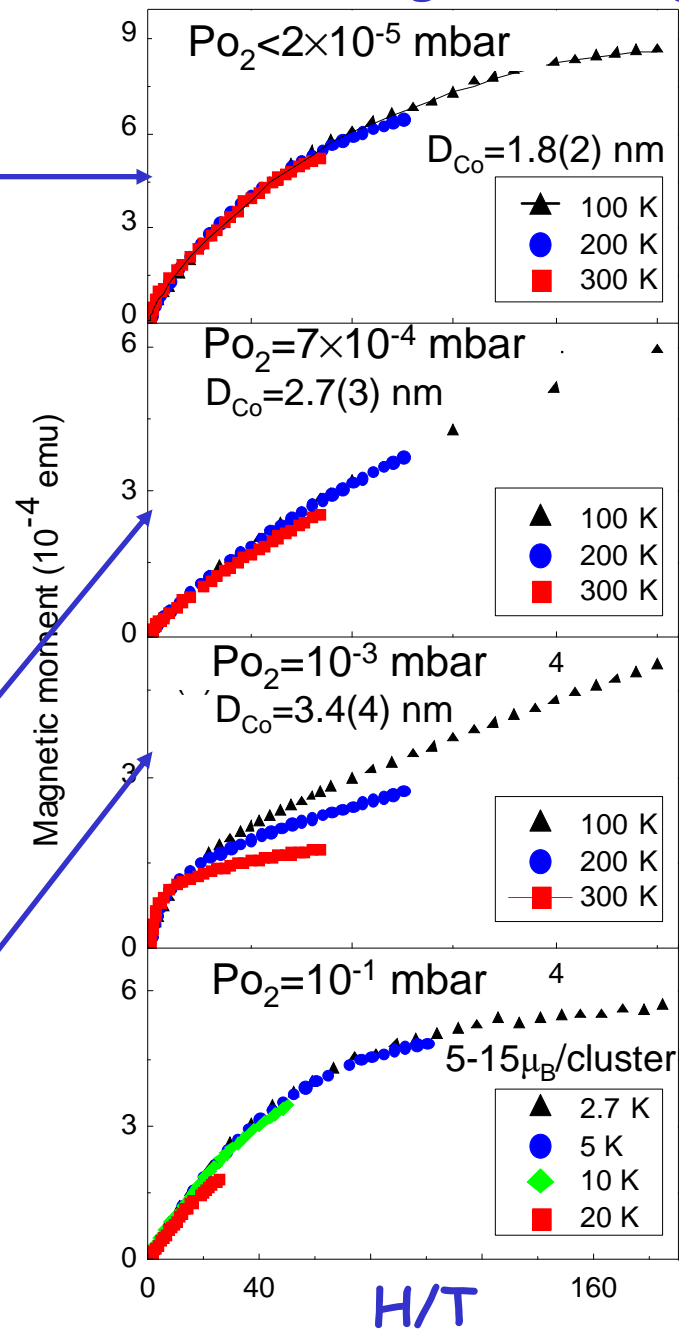


$P_{O_2} = 5 \times 10^{-4}$  mbar

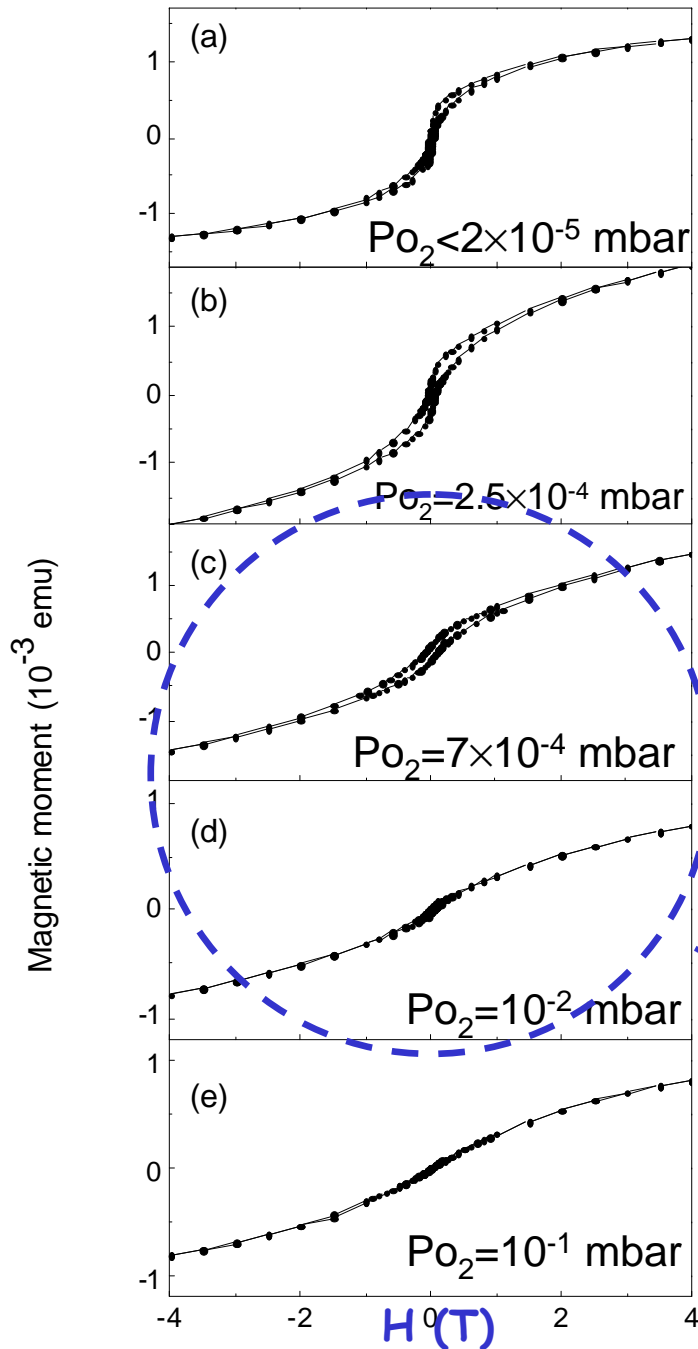
### ZFC-FC curves



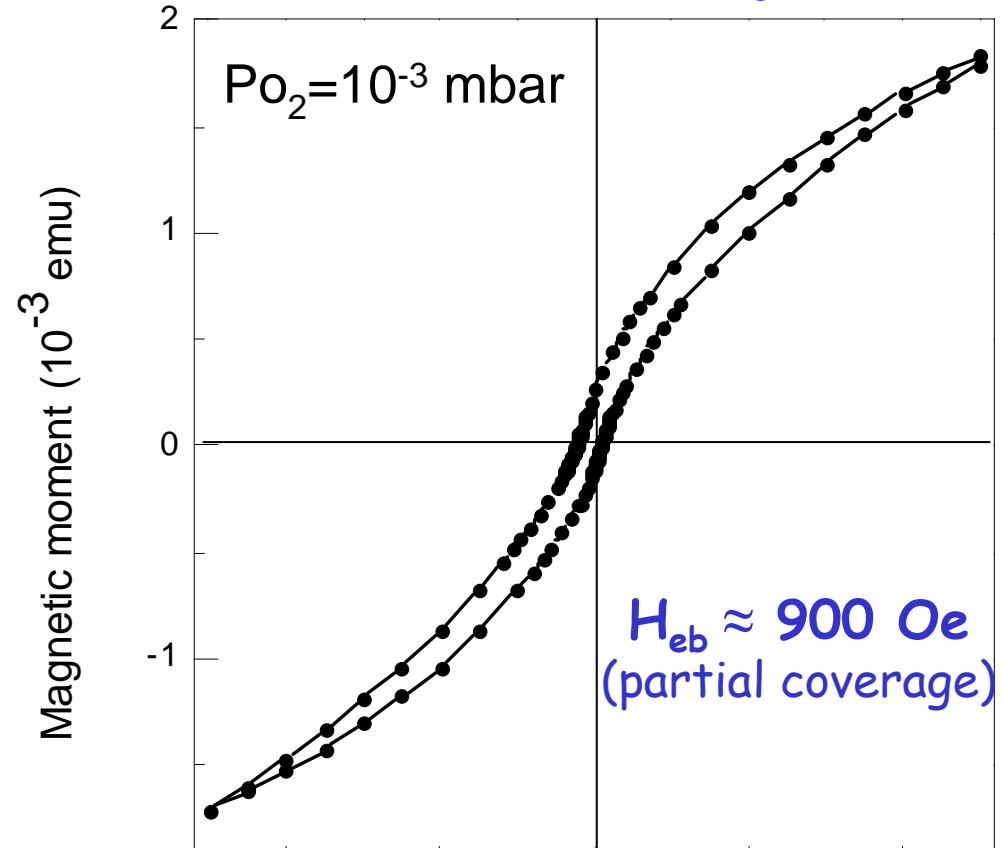
### M vs H/T scaling (SPM region)



# M(H) @ 5 K after ZFC



# M(H) @ 1.8 K after $H_{FC} = 50$ kOe

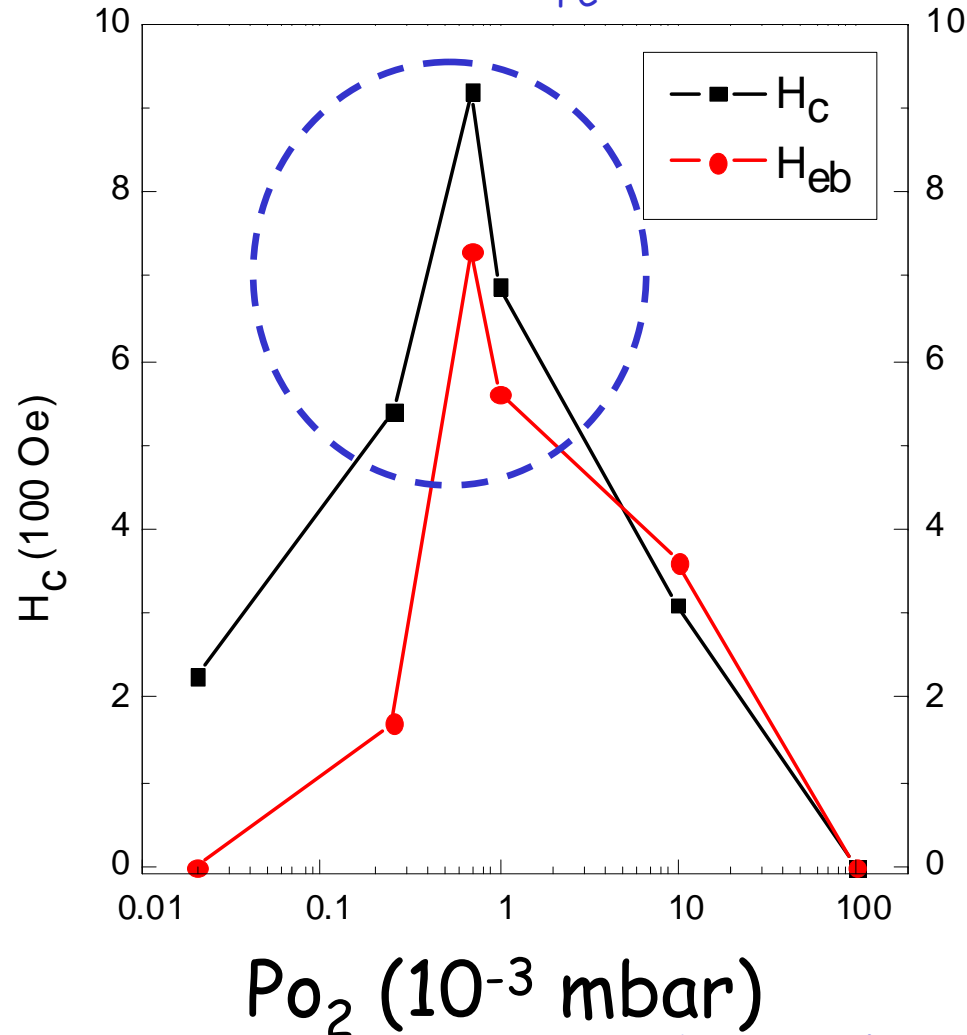


High irreversibility disappears and the small moment in the AF unblocks due to the absence of exchange coupling and magnetic frustration.

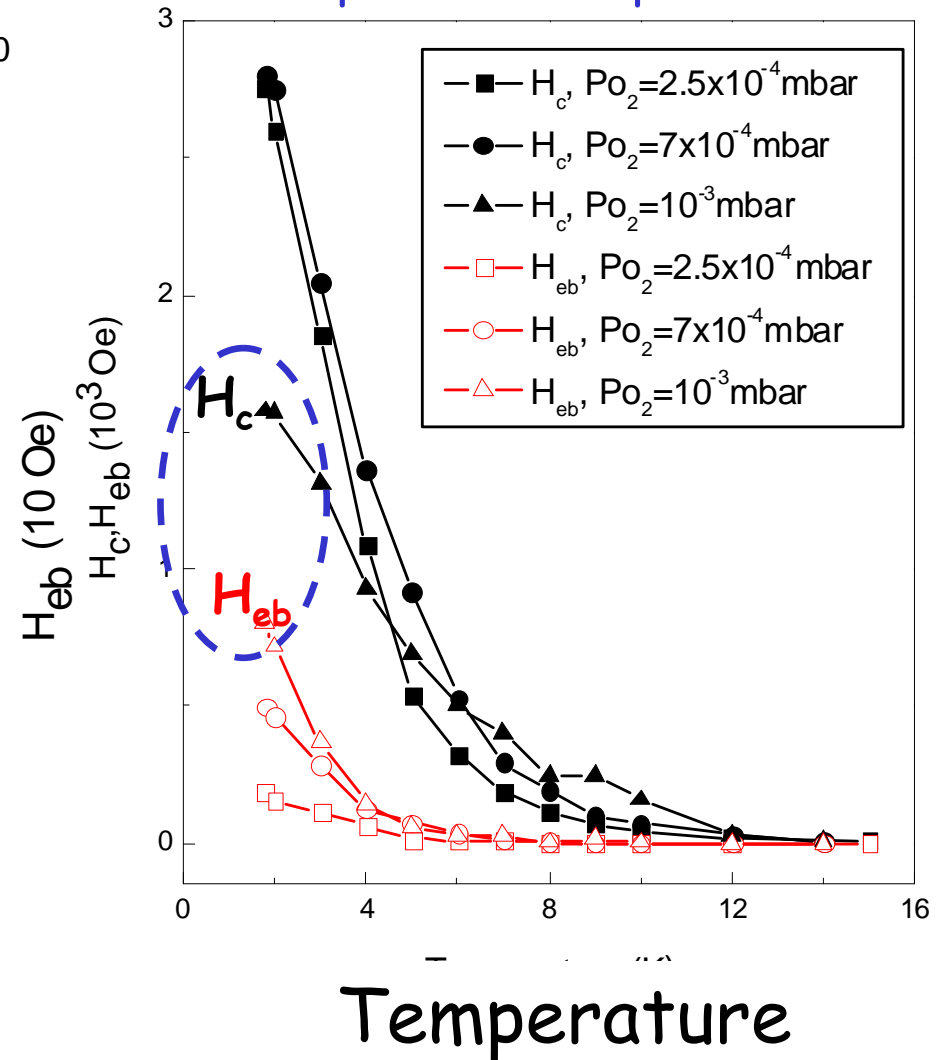
**EB requires a minimum size for both the FM and AF components to show up.**



Coercive field ( $H_c$ ) and  
loop shift ( $H_{eb}$ )  
@ 5 K after  $H_{FC} = 50$  kOe



Temperature dependence



Thermal stabilization of the FM  
due to exchange coupling?




# V. Results:

## Monte Carlo simulations

### References

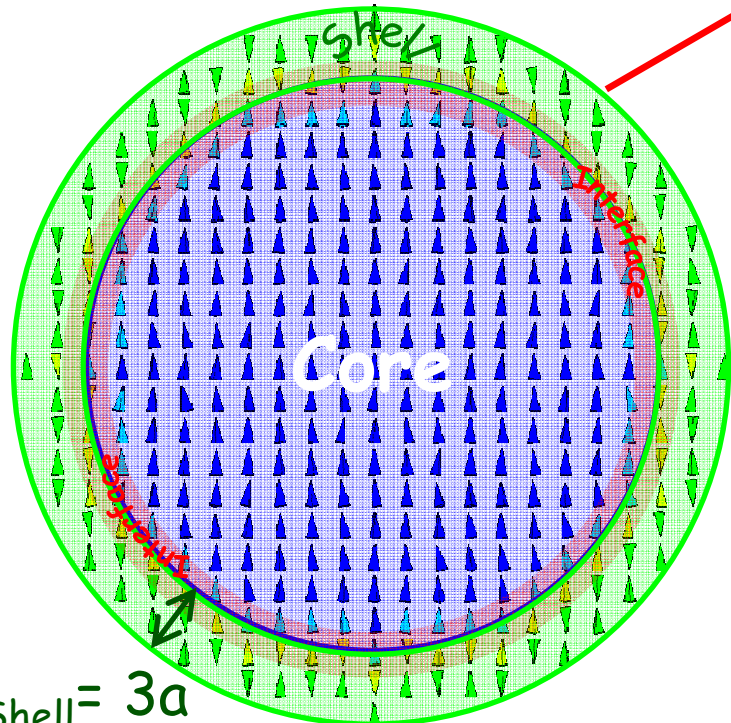
- O. Iglesias, X. Batlle and A. Labarta,  
Phys. Rev. B **72**, 212401 (2005)  
J. Magn. Magn. Mater. **316**, 140-142 (2007)  
J. Phys.: Condens. Matter **19**, 406232 (2007)  
J. Phys. D: Appl. Phys **41**, 134010 (2008)

Exchange bias phenomenology and models of core/shell nanoparticles  
O. Iglesias, A. Labarta and X. Batlle  
Journal of Nanosciences and Nanotechnology **8**, 2761 (2008) REVIEW



Òscar Iglesias  
18:45 (Today)

# Model: single core/shell NP (to avoid collective effects)



Core: ferromagnetic (Co)  
 Shell: antiferromagnetic (oxide)  
 Interface: spins at C (Sh) with nearest neighbors at the Sh (C)

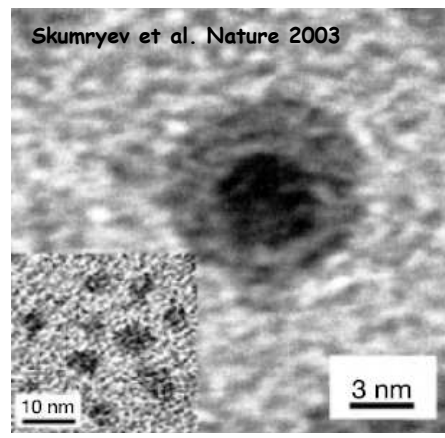
In a core/shell particle, the interface is not well-defined as in bilayers and finite-size effects appear

Interface incorporates roughness, disorder and local compensation/non-compensation (number of neighbors depend on position)

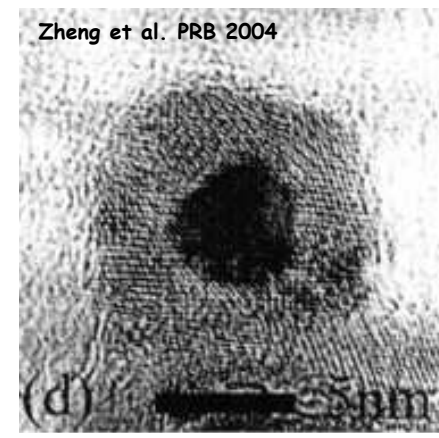
$R_{shell} = 3a$   
 $R_{core} = 9a$   
 $R = 12a$

Number of spins ( $a = 0.3$  nm; unit cell)

$N_{Total} = 5575$   
 $N_{Core} = 3071, N_{Shell} = 2504$   
 $N_{Interface} = 918$  (Shell) +  $794$  (Core)



Co/CoO



Fe/ $\gamma$ -Fe<sub>2</sub>O<sub>3</sub>

$$H/k_B = - \sum_{\langle i,j \rangle} J_{ij} \vec{S}_i \cdot \vec{S}_j - \sum_i K_i (\vec{S}_i \cdot \hat{n}_i)^2 - \vec{h} \cdot \sum_i \vec{S}_i$$

Monte Carlo simulation, Metropolis algorithm for continuous spins  
 $S_i$  = Heisenberg Spins in simple cubic lattice

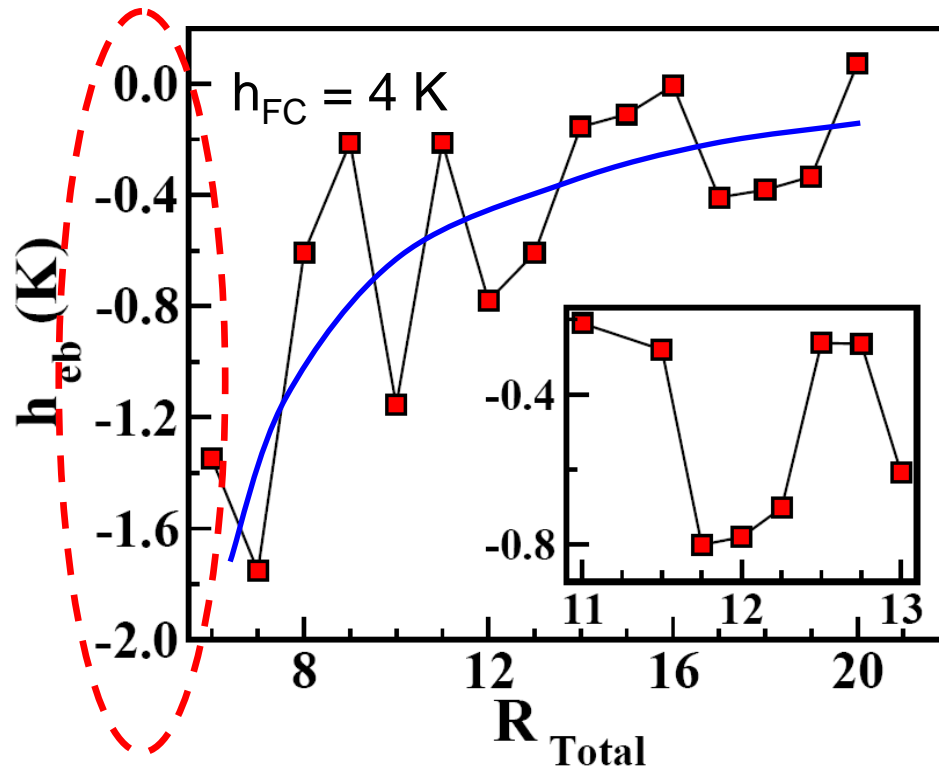
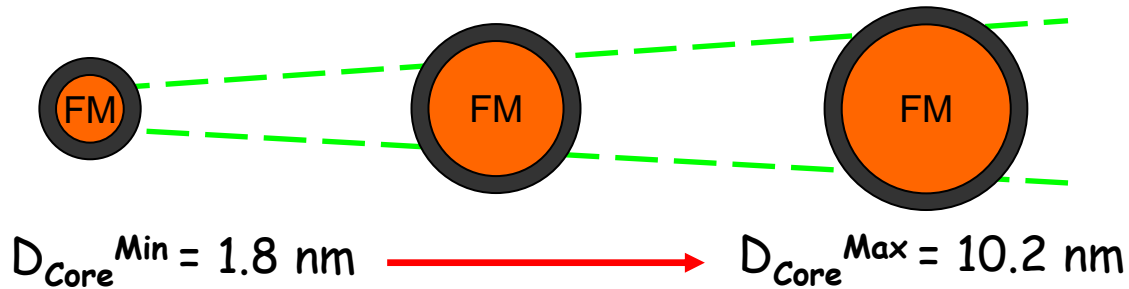
<b>Exchange (n.n.) interaction:</b>
$J_C > 0$ (FM) at the Core
$J_S < 0$ (AF) at the Shell
$J_{Int} > < 0$ (FM or AF) at the Interface
$J_{Int}$ variable within $\pm J_C$

<b>Anisotropy energy</b>
$n_i = z$ axis, uniaxial anisotropy
$K_C$ at the Core
$K_S > K_C$ at the Shell

<b>Zeeman energy</b>
$h$ along $z$ axis
Magnetic field is in temperature units: $h = \mu H/k_B$

Simulation parameters (all in temperature units)			
$J_C = 10$ (fixed)	Fixes Curie temp. $T_C = 29$	$K_C = 1$ (per site)	Fixes coercive field of FM
$J_S = -0.5 J_C$	Fixes Néel temp. $T_N = 14.5 < T_C$	$K_S = 10$ (per site)	Shell with high anisotropy

# Particle size dependence



O. Iglesias, X. Batlle and A. Labarta, J. Phys. D **41**, 134010 (2008)

Oscillatory dependence on particle size

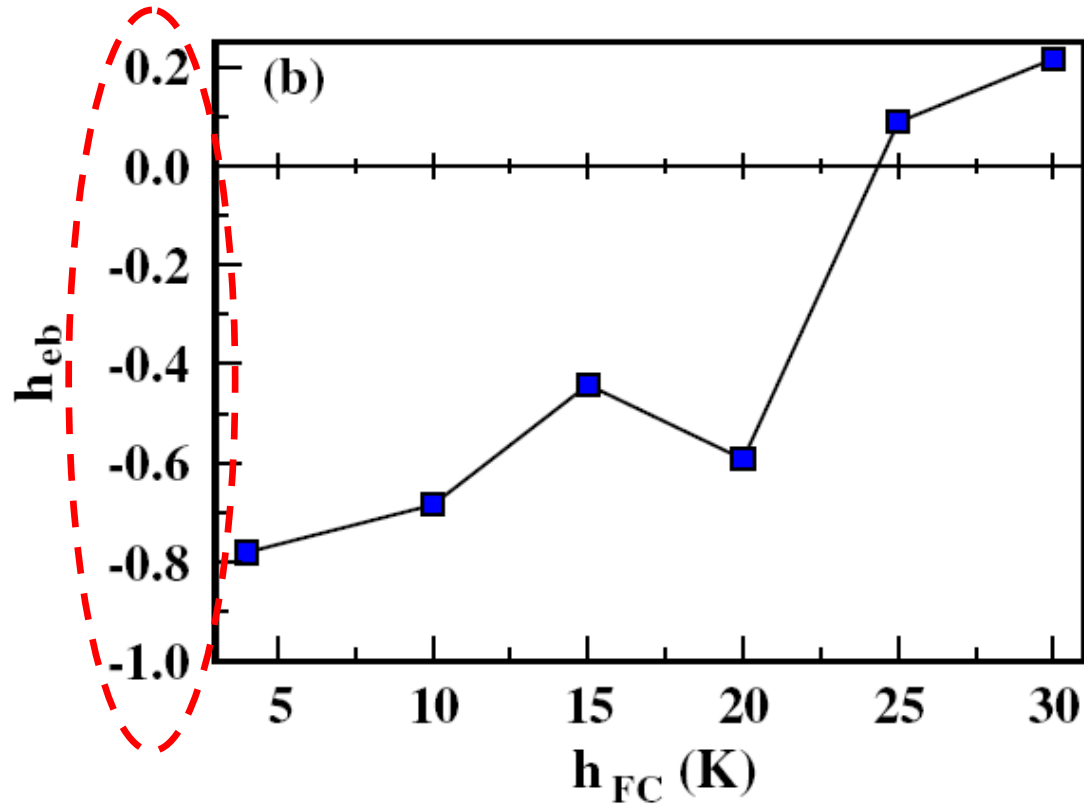
$h_{eb}$  decreases as size increases (experiments)

$$h_{eb} \sim 1/R_{Core}$$

The net magnetization of the AF shell spins at interface oscillates with particle size

Small changes in the core radius induce different geometric arrangements of interfacial shell spins, due to the intersection of the sphere with the lattice sites

# Field cooling dependence



$h_{eb}$  decreases as the cooling field increases (as in experiments)

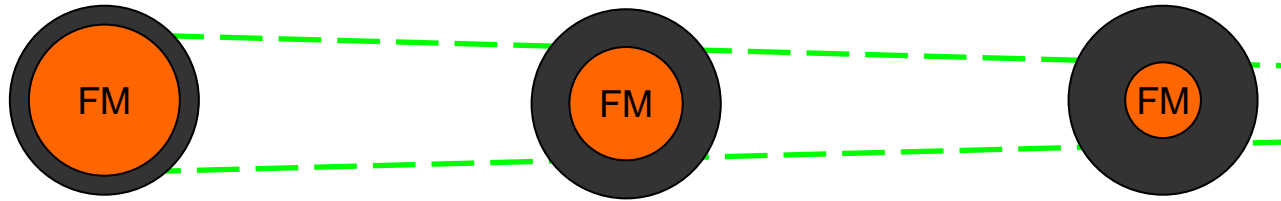
The increasing cooling field progressively reverses the interfacial shell spins along the field direction (as in *positive* EB), reducing the exchange field on the FM

The loops become more symmetric

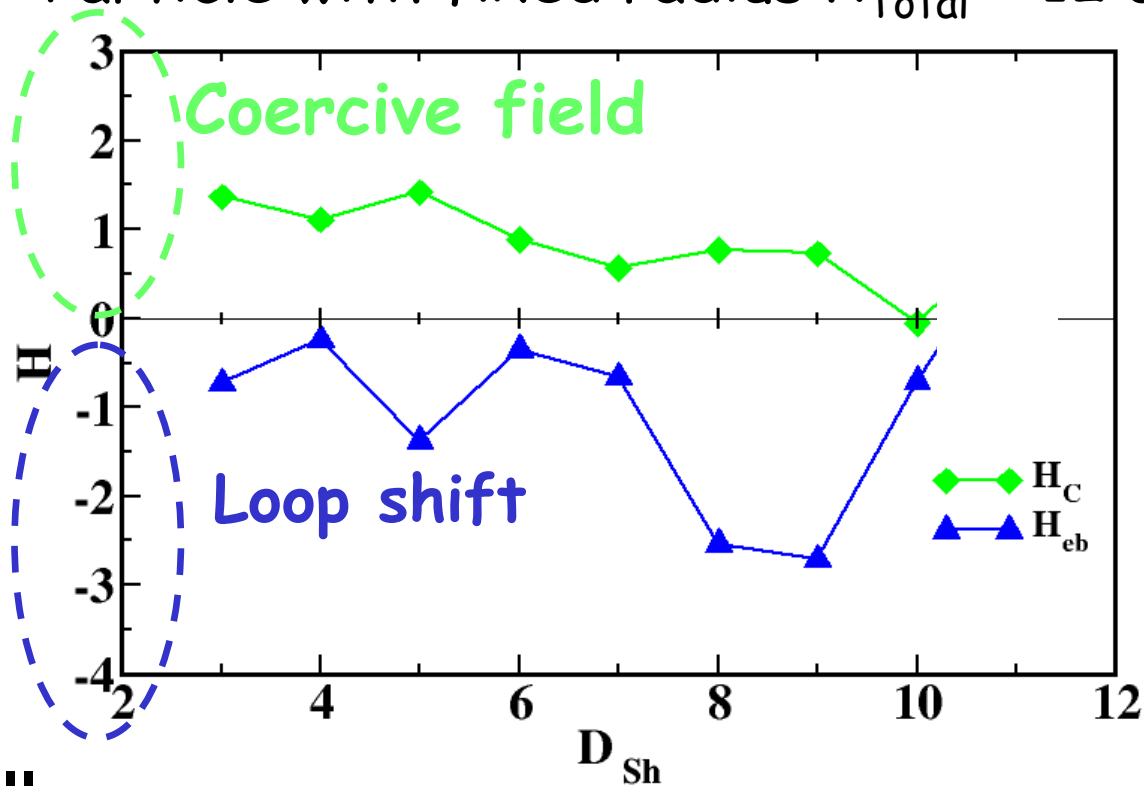
O. Iglesias, X. Batlle and A. Labarta,  
J. Phys. D **41**, 134010 (2008)



# Shell thickness dependence



Particle with fixed radius  $R_{\text{Total}} = 12 \text{ a}$



Core/Shell  
particle



AF  
particle

# Conclusions



1. EB effects and glassy behavior are difficult to decouple.
2. Core/shell NPs naturally incorporate roughness and non-compensation at the interface.
3. EB in core/shell Co-CoO can be tuned as a function of the oxygen pressure.
4. Atomistic Monte Carlo simulations can be used to study finite-size and surface effects at the microscopic level.
5. Simulations of a model of core/shell NP unveils and quantifies the **microscopic origin** of EB: loop shift is due to the exchange field acting on the particle core, generated by the net magnetization of **uncompensated, pinned shell spins at the interface**.
6. EB-related phenomenology such as **particle size, shell thickness, shell anisotropy, cooling field and interface coupling dependences**, together with loop **asymmetry, reversal mechanisms and vertical shift**, are successfully accounted for by the results of the simulation.